

LIEB–ROBINSON BOUNDS FOR THE TODA LATTICE

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ABSTRACT. We establish locality estimates, known as Lieb–Robinson bounds, for the Toda lattice. In contrast to harmonic models, the Lieb–Robinson velocity for these systems do depend on the initial condition. Our results also apply to the entire Toda as well as the Kac–van Moerbeke hierarchy. Under suitable assumptions, our methods also yield a finite velocity for certain perturbations of these systems.

1. INTRODUCTION

Analyzing the dynamics of certain non-relativistic systems is crucial to understanding a number of important problems in statistical mechanics. For example, there is much interest in rigorously justifying the emergence of macroscopic non-equilibrium phenomena, like Fourier’s Law or other forms of heat conduction, directly from a many-body Hamiltonian dynamics [4]. Future progress on fundamental questions such as these will require detailed information on the structure of the underlying dynamics. One of the goals of the present work is to investigate an approximate form of locality, often described by a Lieb–Robinson bound, for the dynamics corresponding to a class of integrable systems.

Recognizing approximate forms of locality, despite the lack of a relativistic framework, has been essential in solving many intriguing open problems. One of the first mathematical formulations of a useful locality estimate was given by Lieb and Robinson in 1972 [24]. In this work, they demonstrated that the dynamics corresponding to quantum spin systems, with e.g. finite-range interactions, remains effectively confined to a ”light” cone, up to corrections which decay at least exponentially away from the light cone. Recently there have been a number of improvements and generalizations of the original result [31, 20, 28, 10, 2, 29, 37, 36, 39, 34], and these new techniques have led to some interesting applications [17, 18, 32, 30, 21, 7, 6, 3, 5, 16]. For a review of these results, we refer the interested reader to [33, 19].

Shortly after the original result of Lieb and Robinson, it was shown in [26], see also [9, 38] for more recent developments, that this notion of quasi-locality also applies to the dynamics of classical oscillator systems. Since it is this work that more closely pertains to the topic of the present article, we will discuss it briefly as follows.

Consider a system of particles confined to a large but finite set $\Lambda \subset \mathbb{Z}^d$. To each site $x \in \Lambda$ associate a particle, or oscillator, with position $q_x \in \mathbb{R}$ and momentum $p_x \in \mathbb{R}$. The state of the system in Λ is described by a sequence $\mathbf{x} = \{(q_x, p_x)\}_{x \in \Lambda}$, and the set of all such sequences, \mathcal{X}_Λ , is called phase space. A Hamiltonian, H , is a real-valued function on phase space. Given a Hamiltonian and a sequence $\mathbf{x} \in \mathcal{X}_\Lambda$, Hamilton’s equations of motion are: for each $x \in \Lambda$,

$$(1.1) \quad \dot{q}_x(t) = \frac{\partial H}{\partial p_x} \quad \text{and} \quad \dot{p}_x(t) = -\frac{\partial H}{\partial q_x}$$

solved with initial condition $\{(q_x(0), p_x(0))\}_{x \in \Lambda} = \mathbf{x}$. For many Hamiltonians, this system of coupled differential equations can be solved for all time (and any initial condition). In this case, we denote by Φ_t , the Hamiltonian flow, i.e., the mapping that associates to initial conditions, the solution at time t : $\Phi_t(\mathbf{x}) = \{(q_x(t), p_x(t))\}_{x \in \Lambda}$.

Measurements of the system under consideration correspond to observables, where an observable A is a complex-valued function on phase space. For example, the position of the particle at site $x \in \Lambda$ corresponds to an observable Q_x for which $Q_x(\mathbf{x}) = q_x$. Let us denote by \mathcal{A}_Λ the set of all observables over \mathcal{X}_Λ . Given

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a Hamiltonian for which (1.1) can be solved, the Hamiltonian dynamics $\alpha_t : \mathcal{A}_\Lambda \rightarrow \mathcal{A}_\Lambda$ is defined by $\alpha_t(A) = A \circ \Phi_t$. In this case, we can make time-dependent observations such as $[\alpha_t(Q_x)](x) = q_x(t)$.

It is well known that the Hamiltonian dynamics is generated by the Poisson bracket, i.e.,

$$(1.2) \quad \frac{d}{dt}\alpha_t(A) = \alpha_t(\{A, H\}) = \{\alpha_t(A), H\} ,$$

and here the Poisson bracket of two observables is the observable

$$(1.3) \quad \{A, B\} = \sum_{x \in \Lambda} \frac{\partial A}{\partial q_x} \cdot \frac{\partial B}{\partial p_x} - \frac{\partial A}{\partial p_x} \cdot \frac{\partial B}{\partial q_x} .$$

In terms of this Poisson bracket, we can now describe the Lieb–Robinson bound.

Let us fix finite sets $X, Y \subset \mathbb{Z}^d$ with $X \cap Y = \emptyset$. Take $\Lambda \subset \mathbb{Z}^d$ finite, but large enough so that $X \cup Y \subset \Lambda$. Consider two observables $A, B \in \mathcal{A}_\Lambda$ with supports in X, Y respectively, i.e., e.g. A depends only on those q_x and p_x with $x \in X$. It is clear that $\{A, B\} = 0$ since A and B have disjoint supports. A Hamiltonian H satisfies a Lieb–Robinson bound if for some initial condition x there exist numbers μ , C , and v for which

$$(1.4) \quad |[\{\alpha_t(A), B\}](x)| \leq C e^{-\mu(d(X, Y) - v|t|)} ,$$

where $d(X, Y)$ is the distance between X and Y . In words, this estimate shows that for times t with $v|t| \leq d(X, Y)$ the support of $\alpha_t(A)$ remains essentially disjoint from the support of B , up to exponentially small corrections in $d(X, Y)$. The number v is often called the Lieb–Robinson velocity corresponding to H , and it represents a bound on the rate at which disturbances can propagate through the system.

We will now briefly discuss what is known concerning Lieb–Robinson bounds for classical systems. The works in [26, 9, 38] prove Lieb–Robinson type bounds for a variety of harmonic systems, and it is shown that analogous results also hold for certain anharmonic perturbations. The strongest such result is in [9] and demonstrates that for an anharmonic model with a quartic on-site perturbation the relevant Poisson bracket decays to zero whenever the distance between the supports of the local observables grows faster than $t \log^\alpha(t)$ for suitable $\alpha > 0$. In [38], where explicit estimates on the Lieb–Robinson velocity were obtained, it is shown that the number v is independent of the initial condition x . Since the system is linear, this is not so surprising.

Of course when it comes to anharmonic lattice oscillations one of the central objects is the famous Fermi–Pasta–Ulam–Tsingou (FPU) problem. It was first demonstrated by Zabusky and Kruskal [44] that one key ingredient to resolve the FPU paradox is the relation with solitons. Moreover, subsequently Toda [43] presented an anharmonic lattice which possess soliton solutions and was later on shown to be integrable by Flaschka [12, 13] and Manakov [25]. Clearly this naturally raises the question about Lieb–Robinson type locality bounds for the Toda lattice and the main goal of this work is to establish such bounds. Our estimates produce a Lieb–Robinson velocity that depends on the initial condition, see e.g. Theorem 2.3. We discuss this fact in the context of one-soliton solutions in Section 2.4. For a restricted class of initial conditions, we can prove the existence of a finite Lieb–Robinson velocity for a class of perturbations of the Toda lattice. We have two bounds of this type. The first, in Section 3.1, establishes a result by directly mimicking the methods for the unperturbed system. The next, in Section 3.2, uses interpolation and shows that the velocity of the perturbed system can be estimated in terms of the velocity of the unperturbed system. Analogous results are shown to also hold for solutions of the Toda hierarchy, see Section 4 and Section 5. For a much larger class of initial conditions, we can prove a locality bound for perturbed systems, however, we do not establish the existence of a finite Lieb–Robinson velocity. This is discussed in Section 6. The final section, Appendix A, describes a set of perturbations and corresponding initial conditions for which solutions remain globally bounded. This shows that the results obtained in sections 3.1 and 3.2 are not vacuous.

2. LOCALITY ESTIMATES FOR THE TODA LATTICE

In this section, we discuss the Toda Lattice and prove a Lieb–Robinson bound. We introduce the model in Section 2.1. The crucial solution estimate, Theorem 2.1, is contained in Section 2.2 as well as some useful remarks. A Lieb–Robinson bound, see Theorem 2.3, in terms of a large class of observables is proven in Section 2.3. We end this section by comparing the velocity estimates in our Lieb–Robinson bound to known results for one-soliton solutions of the Toda Lattice. This is done in Section 2.4.

2.1. The Toda Lattice. The Toda Lattice is a well-studied physical model and one of the prototypical discrete integrable wave equations. We refer to the monographs [11], [40], [43] or the review articles [22], [41] for further information. Like those discussed in the introduction, the Toda Lattice corresponds to a specific system of coupled oscillators on \mathbb{Z} . We present this model immediately in the infinite volume setting.

For each site $n \in \mathbb{Z}$, we associate an ordered pair $(q_n, p_n) \in \mathbb{R}^2$ and denote by \mathcal{X} the set of all sequences $\mathbf{x} = \{(q_n, p_n)\}_{n \in \mathbb{Z}}$. The system is said to be comprised of an infinite collection of *oscillators*, each situated at a site $n \in \mathbb{Z}$ with position $q_n \in \mathbb{R}$ and momentum $p_n \in \mathbb{R}$. Here the state of the system is described by a sequence $\mathbf{x} = \{(q_n, p_n)\}_{n \in \mathbb{Z}} \in \mathcal{X}$.

The oscillators evolve in time according to the following coupled system of differential equations. For each $n \in \mathbb{Z}$ and any $t \in \mathbb{R}$,

$$(2.1) \quad \dot{q}_n(t) = p_n(t) \quad \text{and} \quad \dot{p}_n(t) = e^{-(q_n(t) - q_{n-1}(t))} - e^{-(q_{n+1}(t) - q_n(t))},$$

with initial condition given by some state of the system $\mathbf{x} \in \mathcal{X}$. We note that the differential system (2.1) corresponds to the following (formal) Hamiltonian $H : \mathcal{X} \rightarrow \mathbb{R} \cup \{\infty\}$ given by

$$(2.2) \quad H(\mathbf{x}) = \sum_{n \in \mathbb{Z}} \frac{p_n^2}{2} + V(q_{n+1} - q_n),$$

where $V(r) = e^{-r} + r - 1$. For many sequences, the formal Hamiltonian may be infinite, however, it does formally generate the system of differential equations in (2.1) through Hamilton's equations. For classes of solutions with prescribed decay properties we refer to [42].

Existence and uniqueness of global solutions to (2.1) on certain subsets of \mathcal{X} is well-known. Rather than address this directly, we begin by changing variables. Suppose that (2.1) has a solution. For each $n \in \mathbb{Z}$ and $t \in \mathbb{R}$, set

$$(2.3) \quad a_n(t) = \frac{1}{2} e^{-(q_{n+1}(t) - q_n(t))/2} \quad \text{and} \quad b_n(t) = -\frac{1}{2} p_n(t),$$

in terms of this given solution. This choice is commonly referred to as Flaschka variables as they were initially introduced in [12], see also [13]. Using (2.1), it is clear that these new variables satisfy the following equations of motion

$$(2.4) \quad \dot{a}_n(t) = a_n(t) (b_{n+1}(t) - b_n(t)) \quad \text{and} \quad \dot{b}_n(t) = 2(a_n^2(t) - a_{n-1}^2(t)),$$

and correspond to the following formal Hamiltonian

$$(2.5) \quad H(\mathbf{x}) = \sum_{n \in \mathbb{Z}} 2b_n^2 + V(-\ln(4a_n^2)) = \sum_{n \in \mathbb{Z}} 2b_n^2 + 4a_n^2 - 2\ln(2a_n) - 1.$$

In this work, we will concern ourselves mainly with the properties of the solutions of (2.4).

Consider the vector space $M = \ell^\infty(\mathbb{Z}, \mathbb{R}) \times \ell^\infty(\mathbb{Z}, \mathbb{R})$ of pairs of bounded, real-valued sequences. We will write each $\mathbf{x} \in M$ as

$$(2.6) \quad \mathbf{x} = (\{a_n\}_{n \in \mathbb{Z}}, \{b_n\}_{n \in \mathbb{Z}}) = \{(a_n, b_n)\}_{n \in \mathbb{Z}}.$$

M is a Banach space with respect to the norm

$$(2.7) \quad \|\mathbf{x}\|_M = \max(\|a\|_\infty, \|b\|_\infty), \quad \|c\|_\infty = \sup_n |c_n|.$$

Given any initial condition $\mathbf{x} \in M$, global solvability of the system (2.4) is well-known (see e.g. [40], Theorem 12.6). We will denote the *Toda flow* on M by Φ_t , i.e., $\Phi_t : M \rightarrow M$ is the function that associates an initial condition $\mathbf{x} \in M$ to the solution of (2.4) at time t : $\Phi_t(\mathbf{x}) = \{(a_n(t), b_n(t))\}_{n \in \mathbb{Z}}$ with $\Phi_0(\mathbf{x}) = \mathbf{x}$.

2.2. Estimating Toda solutions. More is known about the solutions of (2.4) on M . In fact, let us fix $\mathbf{x} \in M$ and introduce operators $L(\mathbf{x}), P(\mathbf{x}) : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$ by setting

$$(2.8) \quad [L(\mathbf{x})f]_n = a_n f_{n+1} + a_{n-1} f_{n-1} + b_n f_n$$

and

$$(2.9) \quad [P(\mathbf{x})f]_n = a_n f_{n+1} - a_{n-1} f_{n-1}.$$

For each $\mathbf{x} \in M$, the existence of global solutions imply that $L(\Phi_t(\mathbf{x}))$ and $P(\Phi_t(\mathbf{x}))$ are well-defined for all $t \in \mathbb{R}$. Whenever the initial condition \mathbf{x} is fixed, we will write $L(t) = L(\Phi_t(\mathbf{x}))$ and $P(t) = P(\Phi_t(\mathbf{x}))$ to spare

notation. Observe that for each $x \in M$, $\Phi_t(x) \in M$ and so $L(t)$ is a bounded self-adjoint operator with the operator norm satisfying

$$(2.10) \quad \max(\|a(t)\|_\infty, \|b(t)\|_\infty) \leq \|L(t)\|_2 \leq 2\|a(t)\|_\infty + \|b(t)\|_\infty.$$

Similarly $P(t)$ is a bounded skew-adjoint operator. Since $P(t)$ is also differentiable with respect to t , it generates a two-parameter family of unitary propagators $U(t, s)$ satisfying

$$(2.11) \quad \frac{d}{dt}U(t, s) = P(t)U(t, s) \text{ with } U(t, t) = \mathbb{1},$$

for all pairs $t, s \in \mathbb{R}$, see e.g. [40], Theorem 12.4 for more details. Moreover, a short calculation shows that $P(t)$ and $L(t)$ are a Lax pair associated to (2.4), i.e.,

$$(2.12) \quad \frac{d}{dt}L(t) = [P(t), L(t)],$$

and therefore,

$$L(t) = U(t, s)L(s)U(t, s)^{-1}.$$

From this fact, each $x \in M$ satisfies the a-priori estimate

$$(2.13) \quad \|\Phi_t(x)\|_M \leq \|L(t)\|_2 = \|L(0)\|_2 \quad \text{for all } t \in \mathbb{R}.$$

It is important to emphasize the fact that we work with real-valued solutions. In fact, while a local existence result can be established in the complex case, even periodic complex initial conditions can blow up in finite time (see e.g. [14] or [15]).

Our proof of the Lieb–Robinson bound for the Toda Lattice, see Theorem 2.3, makes crucial use of the following estimate describing the sensitivity of solutions to changes in the initial condition.

Theorem 2.1. *Let $x \in M$ and $\mu > 0$. There exists a number $v = v(\mu, x)$ for which given any $n, m \in \mathbb{Z}$, the bound*

$$(2.14) \quad \max \left[\left| \frac{\partial}{\partial z} a_n(t) \right|, \left| \frac{\partial}{\partial z} b_n(t) \right| \right] \leq \frac{8}{\sqrt{17}} e^{-\mu(|n-m|-v|t|)},$$

holds for all $t \in \mathbb{R}$ and each $z \in \{a_m, b_m\}$. In fact, one may take

$$(2.15) \quad v = \left(1 + \sqrt{17}\right) \|L(0)\|_2 \left(e^{\mu+1} + \frac{1}{\mu}\right).$$

Proof. Fix $x \in M$. Global existence of solutions on M guarantees that for each $x \in M$ and $n \in \mathbb{Z}$, the function $F_n : \mathbb{R} \rightarrow \mathbb{R}^2$ given by

$$(2.16) \quad F_n(t; x) = \begin{pmatrix} a_n(t) \\ b_n(t) \end{pmatrix}$$

is well-defined. It is differentiable with respect to each $z \in \{a_m, b_m\}$, e.g. as a consequence of Lemma 4.1.9 in [1]. When convenient, we will suppress the dependence of F_n on x . Using the equations of motion, i.e. (2.4), it is clear that

$$(2.17) \quad F_n(t) = F_n(0) + \int_0^t \begin{pmatrix} a_n(s)(b_{n+1}(s) - b_n(s)) \\ 2(a_n^2(s) - a_{n-1}^2(s)) \end{pmatrix} ds.$$

The relation

$$(2.18) \quad \frac{\partial}{\partial z} F_n(t) = \frac{\partial}{\partial z} F_n(0) + \sum_{|e| \leq 1} \int_0^t D_{n,e}(s) \frac{\partial}{\partial z} F_{n+e}(s) ds,$$

with

$$(2.19) \quad D_{n,e}(s) = \begin{pmatrix} (b_{n+1}(s) - b_n(s)) \delta_0(e) & a_n(s)(-\delta_0(e) + \delta_1(e)) \\ 4(a_n(s)\delta_0(e) - a_{n-1}(s)\delta_{-1}(e)) & 0 \end{pmatrix}.$$

then follows immediately from (2.17).

To complete our estimate, we introduce the following notation. For each $v \in \mathbb{R}^2$, we will denote by

$$(2.20) \quad v = \begin{pmatrix} x \\ y \end{pmatrix} \quad \text{and} \quad |v| = \begin{pmatrix} |x| \\ |y| \end{pmatrix}.$$

Moreover, we will write

$$(2.21) \quad \left| \begin{pmatrix} x \\ y \end{pmatrix} \right| \leq \left| \begin{pmatrix} u \\ v \end{pmatrix} \right| \quad \text{if and only if} \quad |x| \leq |u| \quad \text{and} \quad |y| \leq |v|.$$

With this understanding, the uniform solution estimate (2.13), and (2.19), it is clear that

$$(2.22) \quad \left| \frac{\partial}{\partial z} F_n(t) \right| \leq \left| \frac{\partial}{\partial z} F_n(0) \right| + \|L(0)\|_2 \sum_{|e| \leq 1} \int_0^{|t|} D_e \left| \frac{\partial}{\partial z} F_{n+e}(s) \right| ds,$$

where

$$(2.23) \quad D_e = \begin{pmatrix} 2\delta_0(e) & \delta_0(e) + \delta_1(e) \\ 4(\delta_0(e) + \delta_{-1}(e)) & 0 \end{pmatrix}.$$

Let us now consider the case that $z = a_m$, i.e.,

$$(2.24) \quad \frac{\partial}{\partial a_m} F_n(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \delta_m(n).$$

In this case, iteration of (2.22) yields

$$(2.25) \quad \begin{aligned} \left| \frac{\partial}{\partial a_m} F_n(t) \right| &\leq \sum_{k=0}^{\infty} \frac{(\|L(0)\|_2 |t|)^k}{k!} \sum_{|e_1| \leq 1} \cdots \sum_{|e_k| \leq 1} \delta_{m+e_1+\dots+e_k}(n) D_{e_1} \cdots D_{e_k} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ &\leq \sum_{k=|n-m|}^{\infty} \frac{(\|L(0)\|_2 |t|)^k}{k!} D^k \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \end{aligned}$$

where we have set

$$(2.26) \quad D = \sum_{|e| \leq 1} D_e = \begin{pmatrix} 2 & 2 \\ 8 & 0 \end{pmatrix}.$$

Moreover, note that the remainder term in finite iterations of (2.22) converges to zero since $\frac{\partial}{\partial z} F_n(t)$ is continuous and thus bounded on compact time intervals.

The eigenvalues of D are $\lambda_{\pm} = 1 \pm \sqrt{17}$, and in terms of the eigenvectors v_{\pm} given by

$$(2.27) \quad v_{\pm} = \begin{pmatrix} \lambda_{\pm} \\ 8 \end{pmatrix} \quad \text{it is clear that} \quad \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{2\sqrt{17}} v_+ - \frac{1}{2\sqrt{17}} v_-.$$

Taking the infinity norm, in \mathbb{R}^2 , of both sides of (2.25) shows that

$$(2.28) \quad \left\| \frac{\partial}{\partial a_m} F_n(t) \right\|_{\infty} \leq \frac{8}{\sqrt{17}} \sum_{k=|n-m|}^{\infty} \frac{(\lambda_+ \|L(0)\|_2 |t|)^k}{k!}$$

Now, let $\mu > 0$ be fixed and set $c = \lambda_+ \|L(0)\|_2$. If $c|t| \leq |n-m|e^{-(\mu+1)}$, then by Stirling

$$(2.29) \quad \sum_{k=|n-m|}^{\infty} \frac{(c|t|)^k}{k!} \leq \frac{(c|t|)^{|n-m|}}{(n-m)!} e^{c|t|} \leq \left(\frac{c|t|}{|n-m|} \right)^{|n-m|} e^{|n-m|} e^{c|t|} \leq e^{-\mu(|n-m| - \frac{c}{\mu}|t|)}.$$

Otherwise $0 \leq -\mu(|n-m| - ce^{\mu+1}|t|)$, and so by (2.28) we have

$$(2.30) \quad \left\| \frac{\partial}{\partial a_m} F_n(t) \right\|_{\infty} \leq \frac{8}{\sqrt{17}} e^{c|t|} \leq \frac{8}{\sqrt{17}} e^{-\mu(|n-m| - c(e^{\mu+1} + \frac{1}{\mu})|t|)}.$$

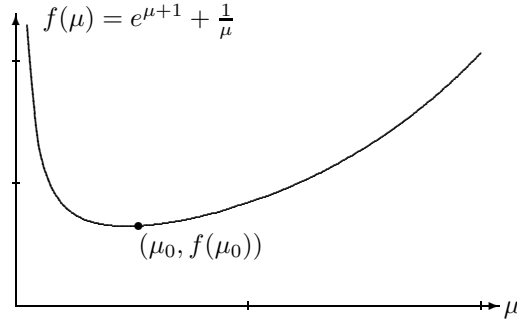
In the case that $z = b_m$, it is clear that

$$(2.31) \quad \begin{pmatrix} 0 \\ 1 \end{pmatrix} = -\frac{\lambda_-}{8(\lambda_+ - \lambda_-)} v_+ + \frac{\lambda_+}{8(\lambda_+ - \lambda_-)} v_-.$$

A short calculation leads to a prefactor of $\frac{\lambda_+}{\sqrt{17}}$ in the analogue of (2.28), and this is less than the one above. We have proven the result. \square

Remarks:

1. The number v in our estimate (2.14) depends on the initial condition x only through the quantity $\|L(0)\|_2$. In fact, one can replace $\|L(0)\|_2$ with any uniform estimate on the solutions (uniform in n and t) as is clear in the bound from (2.18) to (2.22).
2. If $|n - m| > 0$, the left hand side of (2.14) is zero when $t = 0$. The estimate in (2.14) then shows that the solutions of the Toda Lattice at site n are insensitive (i.e. exponentially small) to changes in the initial condition at site m for times t satisfying $v|t| \leq |n - m|$. For this reason, the number v is often called a bound on the velocity (i.e. rate) at which disturbances propagate through the Toda Lattice. Moreover, if $|n - m| \geq 1$, then the bound in (2.28) is at most linear in $|t|$, for small t .
3. In general, the fact that solutions of the Toda Lattice have partial derivatives that decay exponentially in $|n - m|$ will be more important than any particular decay rate μ . For this reason, given an initial condition $x \in M$, one can optimize the quantity v over all possible $\mu > 0$. The graph of $f(\mu) = e^{\mu+1} + \frac{1}{\mu}$ is given below.



It is clear that the optimal μ is achieved when

$$(2.32) \quad \frac{d}{d\mu} \left(e^{\mu+1} + \frac{1}{\mu} \right) = 0$$

and moreover, this value is independent of the initial condition x . The unique solution of (2.32) can be expressed in terms of the Lambert W -function ([35, §4.13]) and is given by

$$(2.33) \quad \mu_0 = 2W(1/(2\sqrt{e})) \approx 0.47767 \quad \text{and} \quad f(\mu_0) \approx 6.47622.$$

2.3. A Lieb–Robinson bound for Toda. We can now formulate a Lieb–Robinson bound for the Toda Lattice. In general, an *observable* A is a function $A : M \rightarrow \mathbb{C}$. Let us denote by \mathcal{A} the set of all observables with well-defined first order partial derivatives. We will say that an observable A is *pointwise bounded* if

$$(2.34) \quad \|A\|_x = \sup_{t \in \mathbb{R}} |A(\Phi_t(x))| < \infty \quad \forall x \in M.$$

Take $\mathcal{A}^{(1)}$ to be the set of those observables in \mathcal{A} for which all of the first order partial derivatives are bounded and

$$(2.35) \quad \|A\|_{1,x} = \sum_{n \in \mathbb{Z}} \left(\left\| \frac{\partial A}{\partial a_n} \right\|_x + \left\| \frac{\partial A}{\partial b_n} \right\|_x \right) < \infty \quad \forall x \in M.$$

An observable $A \in \mathcal{A}$ is said to be *supported* on $X \subset \mathbb{Z}$ if the observables $\frac{\partial A}{\partial a_n}$ and $\frac{\partial A}{\partial b_n}$ are identically zero for all $n \in \mathbb{Z} \setminus X$. The *support* of an observable A is the minimal set on which A is supported, and we will denote this set by $\text{supp}(A)$. Set $\mathcal{A}_0 \subset \mathcal{A}$ to be the set of all observables with compact support. For any $t \in \mathbb{R}$ and $A \in \mathcal{A}$, the *Toda dynamics*, which we denote by α_t , is the observable-valued mapping given by

$$(2.36) \quad \alpha_t(A) = A \circ \Phi_t,$$

where Φ_t is the Toda flow described at the end of Section 2.1 above. In other words, for each $x \in M$, $t \in \mathbb{R}$, and $A \in \mathcal{A}$,

$$(2.37) \quad [\alpha_t(A)](x) = A(\Phi_t(x)),$$

which, since solutions are global on M , is a well-defined quantity.

Example 2.2. The two most basic observables correspond to evaluation maps, i.e., to each $n \in \mathbb{Z}$, we associate the functions A_n and B_n given by

$$(2.38) \quad A_n(x) = a_n \quad \text{and} \quad B_n(x) = b_n \quad \text{for each } x \in M.$$

Each of these observables have support $X = \{n\}$. Moreover, the dynamically evolved observables

$$(2.39) \quad [\alpha_t(A_n)](x) = a_n(t) \quad \text{and} \quad [\alpha_t(B_n)](x) = b_n(t)$$

correspond to observations over time.

The *modified Poisson bracket* of two observables A and B is formally defined as the observable

$$(2.40) \quad \{[A, B]\}(x) = \frac{1}{4} \sum_{n \in \mathbb{Z}} a_n \cdot \left[\frac{\partial A}{\partial a_n} \cdot \frac{\partial B}{\partial \tilde{b}_n} - \frac{\partial A}{\partial \tilde{b}_n} \cdot \frac{\partial B}{\partial a_n} \right](x)$$

where we have denote by $\frac{\partial}{\partial b_n} = \frac{\partial}{\partial b_{n+1}} - \frac{\partial}{\partial b_n}$. If, for example, either A or B has compact support, then the corresponding modified Poisson bracket is a well-defined observable. This quantity is of particular interest since it generates the Toda dynamics, i.e., for any $A \in \mathcal{A}_0$

$$(2.41) \quad \frac{d}{dt} \alpha_t(A) = \alpha_t(\{A, H\}) = \{\alpha_t(A), H\},$$

where H is the Hamiltonian (2.5). Now, for any $n, m \in \mathbb{Z}$, $t \in \mathbb{R}$, and $x \in M$, it is easy to see that

$$(2.42) \quad \{\alpha_t(A_n), B_m\}(x) = \frac{1}{4} a_{m-1} \frac{\partial}{\partial a_{m-1}} a_n(t) - \frac{1}{4} a_m \frac{\partial}{\partial a_m} a_n(t)$$

using the basic observables defined in Example 2.2. In this simple case, Theorem 2.1 gives the bound

$$(2.43) \quad \begin{aligned} |\{\alpha_t(A_n), B_m\}(x)| &\leq \frac{2\|a\|_\infty}{\sqrt{17}} \left(e^{-\mu(|n-m+1|-v|t|)} + e^{-\mu(|n-m|-v|t|)} \right) \\ &\leq \frac{2\|a\|_\infty}{\sqrt{17}} (1 + e^\mu) e^{-\mu(|n-m|-v|t|)}. \end{aligned}$$

In general, we have the following Lieb–Robinson bound for the Toda Lattice.

Theorem 2.3. *Let $x \in M$ and $\mu > 0$. There exist numbers C and v for which given any observables $A, B \in \mathcal{A}^{(1)}$, the estimate*

$$(2.44) \quad |\{\alpha_t(A), B\}(x)| \leq C \|a\|_\infty \sum_{n, m \in \mathbb{Z}} \left(\left\| \frac{\partial A}{\partial a_m} \right\|_x + \left\| \frac{\partial A}{\partial b_m} \right\|_x \right) \left(\left\| \frac{\partial B}{\partial a_n} \right\|_x + \left\| \frac{\partial B}{\partial b_n} \right\|_x \right) e^{-\mu(|n-m|-v|t|)}$$

holds for all $t \in \mathbb{R}$. Here $C = \frac{2}{\sqrt{17}}(1 + e^\mu)$ and v is as in Theorem 2.1.

Proof. Substituting into (2.40), we find that

$$(2.45) \quad \{\alpha_t(A), B\}(x) = \frac{1}{4} \sum_{n \in \mathbb{Z}} a_n \cdot \left[\frac{\partial}{\partial a_n} \alpha_t(A) \cdot \frac{\partial B}{\partial \tilde{b}_n} - \frac{\partial}{\partial \tilde{b}_n} \alpha_t(A) \cdot \frac{\partial B}{\partial a_n} \right](x).$$

Using the chain rule, it is clear that for any $z \in \{a_n, \tilde{b}_n\}$,

$$(2.46) \quad \left[\frac{\partial}{\partial z} \alpha_t(A) \right](x) = \sum_{m \in \mathbb{Z}} \left[\alpha_t \left(\frac{\partial A}{\partial a_m} \right) \cdot \frac{\partial}{\partial z} \alpha_t(A_m) + \alpha_t \left(\frac{\partial A}{\partial b_m} \right) \cdot \frac{\partial}{\partial z} \alpha_t(B_m) \right](x)$$

where we are using the notation from Example 2.2. Applying the triangle inequality and then Theorem 2.1, the bound

$$(2.47) \quad \begin{aligned} |\{\alpha_t(A), B\}(x)| &\leq \frac{\|a\|_\infty}{4} \sum_{n, m} \left(\left\| \frac{\partial A}{\partial a_m} \right\|_x \left| \frac{\partial}{\partial a_n} a_m(t) \right| + \left\| \frac{\partial A}{\partial b_m} \right\|_x \left| \frac{\partial}{\partial a_n} b_m(t) \right| \right) \left\| \frac{\partial B}{\partial \tilde{b}_n} \right\|_x \\ &\quad + \frac{\|a\|_\infty}{4} \sum_{n, m} \left(\left\| \frac{\partial A}{\partial a_m} \right\|_x \left| \frac{\partial}{\partial \tilde{b}_n} a_m(t) \right| + \left\| \frac{\partial A}{\partial b_m} \right\|_x \left| \frac{\partial}{\partial \tilde{b}_n} b_m(t) \right| \right) \left\| \frac{\partial B}{\partial a_n} \right\|_x \\ &\leq \frac{2\|a\|_\infty}{\sqrt{17}} \sum_{n, m} \left(\left\| \frac{\partial A}{\partial a_m} \right\|_x + \left\| \frac{\partial A}{\partial b_m} \right\|_x \right) \left\| \frac{\partial B}{\partial \tilde{b}_n} \right\|_x e^{-\mu(|n-m|-v|t|)} \\ &\quad + \frac{2\|a\|_\infty}{\sqrt{17}} e^\mu \sum_{n, m} \left(\left\| \frac{\partial A}{\partial a_m} \right\|_x + \left\| \frac{\partial A}{\partial b_m} \right\|_x \right) \left\| \frac{\partial B}{\partial a_n} \right\|_x e^{-\mu(|n-m|-v|t|)} \end{aligned}$$

readily follows. Note that the bound on e.g. $\left| \frac{\partial}{\partial b_n} a_m(t) \right|$ follows from the argument in Theorem 2.1. Another triangle inequality proves (2.44) as claimed. \square

For many applications, the observables of interest will have disjoint supports. In this case, the bound in (2.44) can be stated as follows. For any $x \in M$ and $\mu > 0$, there exist numbers C and v , as above, such that given any two disjoint subsets X and Y of \mathbb{Z} the bound

$$(2.48) \quad |\{\alpha_t(A), B\}(x)| \leq C \|a\|_\infty \|A\|_{1,x} \|B\|_{1,x} e^{-\mu(d(X,Y)-v|t|)},$$

holds for all $t \in \mathbb{R}$ and any observables $A, B \in \mathcal{A}^{(1)}$ with supports in X, Y , respectively. Here

$$(2.49) \quad d(X, Y) = \inf\{|x - y| : x \in X \text{ and } y \in Y\} > 0,$$

and (2.48) corresponds with the bound (1.4) claimed in the introduction. Other norms on the observables under considerations are often useful. For example, consider

$$(2.50) \quad \|\partial A\|_x = \sup_{m \in \mathbb{Z}} \max \left[\left\| \frac{\partial A}{\partial a_m} \right\|_x, \left\| \frac{\partial A}{\partial b_m} \right\|_x \right]$$

for any $A \in \mathcal{A}^{(1)}$. Theorem 2.3 then shows that: For any $x \in M$ and $\mu > 0$, there exist numbers C and v , as above, such that given any two disjoint subsets X and Y of \mathbb{Z} , at least one of which being finite, the bound

$$(2.51) \quad |\{\alpha_t(A), B\}(x)| \leq 4C \|a\|_\infty \|\partial A\|_x \|\partial B\|_x \sum_{m \in X, n \in Y} e^{-\mu(|m-n|-v|t|)},$$

holds for all $t \in \mathbb{R}$ and any observables $A, B \in \mathcal{A}^{(1)}$ with supports in X, Y , respectively.

In any case, the bound in Theorem 2.3 demonstrates that each $x \in M$ propagates no faster than some finite rate. We say that the Lieb–Robinson velocity corresponding to $x \in M$ is $v(x, \mu_0)$ where μ_0 is the infimum defined in the remarks after Theorem 2.1.

2.4. On one-soliton solutions. Perhaps the most important collection of solutions to the Toda Lattice are the solitons, i.e. the solitary waves, see e.g. [40, 41, 43]. In fact, solitons can be considered as the stable part of any short-range initial condition since every such solution eventually splits into a number of stable solitons plus a decaying dispersive tail [22]. Since all involved quantities can be computed explicitly for the one-soliton solution we will use it as a test case and for our Lieb–Robinson bounds.

Fix $\kappa > 0$. A one-soliton solution of the Toda lattice is given by

$$(2.52) \quad q_n(t; \pm) = q - \ln \left(\frac{1 + \exp[-2\kappa n \pm 2 \sinh(\kappa)t + \delta]}{1 + \exp[-2\kappa(n-1) \pm 2 \sinh(\kappa)t + \delta]} \right),$$

where q and δ are real constants. Here $q_n(t; \pm)$ represents the position of the traveling wave. It describes a single bump traveling with speed $\pm \frac{\sinh(\kappa)}{\kappa}$ and width proportional to $1/\kappa$. In other words, the smaller the soliton the faster it propagates. Changing δ amounts to a shift of the solution and we will set $\delta = 0$ for simplicity. From this explicit formula it is clear that, in contrast to harmonic models, the velocity of a solution to the Toda Lattice may indeed depend on the initial condition.

In terms of the function

$$(2.53) \quad f_\pm(x, t) = 1 + \exp[-2\kappa x \pm 2 \sinh(\kappa)t]$$

it is clear that

$$(2.54) \quad q_n(t; \pm) = q - \ln \left(\frac{f_\pm(n, t)}{f_\pm(n-1, t)} \right),$$

and from Hamilton's equations, we also know that

$$(2.55) \quad p_n(t; \pm) = \frac{d}{dt} q_n(t; \pm) = \frac{f'_\pm(n-1, t)}{f_\pm(n-1, t)} - \frac{f'_\pm(n, t)}{f_\pm(n, t)}.$$

In Flaschka's variables, we have that

$$(2.56) \quad a_n(t; \pm) = \frac{1}{2} \frac{\sqrt{f_\pm(n-1, t) f_\pm(n+1, t)}}{f_\pm(n, t)} \quad \text{and} \quad b_n(t; \pm) = \frac{1}{2} \left(\frac{f'_\pm(n, t)}{f_\pm(n, t)} - \frac{f'_\pm(n-1, t)}{f_\pm(n-1, t)} \right).$$

A short calculation shows that

$$(2.57) \quad \sup_n a_n(0; \pm) = a_0(0; \pm) = \frac{\cosh(\kappa)}{2} \quad \text{and} \quad \sup_n |b_n(0; \pm)| = |b_0(0; \pm)| = \frac{\sinh(\kappa) \tanh(\kappa)}{2}.$$

With x_κ denoting the initial condition corresponding to this one-soliton, we have proven that

$$(2.58) \quad \frac{\cosh(\kappa)}{2} = \max \left(\frac{\cosh(\kappa)}{2}, \frac{\sinh(\kappa) \tanh(\kappa)}{2} \right) \leq \|L(x_\kappa)\|_2 \leq \cosh(\kappa) + \frac{\sinh(\kappa) \tanh(\kappa)}{2},$$

using e.g. (2.10). Since $L(x_\kappa)$ is self-adjoint its norm is equal to the spectral radius. Moreover, the spectrum is given by an absolutely continuous part $[-1, 1]$ plus the single eigenvalue $\pm \cosh(\kappa)$ implying

$$(2.59) \quad \|L(x_\kappa)\|_2 = \cosh(\kappa).$$

To see this last claim note that the one-soliton solution can be computed from the inverse scattering transform by choosing one eigenvalue $\lambda = \pm \cosh(\kappa)$ plus the corresponding norming constant $\gamma = (1 - e^{-2\kappa})e^\delta$ and zero reflection coefficient (cf. [43, Sect. 3.6] or [40, Sec. 13.4]) or by using the double commutation method to add one eigenvalue λ with norming constant γ to the trivial solution (cf. [40, Sect. 14.5]).

As is clear from this calculation, the Lieb–Robinson velocity does provide a reasonable estimate on the actual velocity, at least for one-soliton solutions.

3. ESTIMATES ON PERTURBED TODA SYSTEMS

In this section, we consider a class of perturbations of the Toda system for which locality results analogous to Theorem 2.3 still hold. We introduce these perturbations as follows. Let $W : \mathbb{R} \rightarrow [0, \infty)$ satisfy $W \in C^2(\mathbb{R})$. Consider the formal Hamiltonian

$$(3.1) \quad H^w = H + \sum_{n \in \mathbb{Z}} W_n,$$

where H is the Toda Hamiltonian, see (2.5), and for any initial condition $x = \{(a_n, b_n)\}_{n \in \mathbb{Z}} \in M_0 = \ell^\infty(\mathbb{Z}, \mathbb{R} \setminus \{0\}) \times \ell^\infty(\mathbb{Z}, \mathbb{R})$, the observable $W_n(x) = W(\ln(4a_n^2))$. Our choice of parametrization is motivated by the fact that the potential V for the unperturbed Toda Lattice, see (2.2), is a function of $q_{n+1} - q_n = -\ln(4a_n^2)$. Since the choice $a_n = 0$ corresponds to taking the positions of the particles at sites n and $n+1$ infinitely far apart, we will exclude it from our considerations. The formal Hamiltonian (3.1) corresponds to the following system of coupled differential equations

$$(3.2) \quad \begin{aligned} \dot{a}_n^w(t) &= a_n^w(t) (b_{n+1}^w(t) - b_n^w(t)) \\ \dot{b}_n^w(t) &= 2(a_n^w(t)^2 - a_{n-1}^w(t)^2) + R_n(t), \end{aligned}$$

where

$$(3.3) \quad R_n(t) = \frac{1}{2} [W'(\ln(4a_n^w(t)^2)) - W'(\ln(4a_{n-1}^w(t)^2))].$$

Local existence and uniqueness of solutions of (3.2) corresponding to initial conditions $x \in M_0$ follows from standard results, [1, Thm. 4.1.5]. Let us denote by Φ_t^w the perturbed Toda flow, i.e., the function that associates initial conditions $x \in M_0$ with $\Phi_t^w(x) = \{(a_n^w(t), b_n^w(t))\}$, the solution of (3.2) at time t . Consider the set of initial conditions $M_b = M_b(W) \subset M_0$ for which there exists numbers $C_1, C_2 < \infty$ with

$$(3.4) \quad \sup_{t \in \mathbb{R}} \|\Phi_t^w(x)\|_M \leq C_1 \quad \text{and} \quad \sup_{t \in \mathbb{R}} \sup_{n \in \mathbb{Z}} \frac{1}{|a_n^w(t)|} \leq C_2.$$

It will be shown in Appendix A that this M_b contains at least all initial conditions whose energy is finite under appropriate assumptions on W .

For initial conditions $x \in M_b$, we have two estimates on the corresponding Lieb–Robinson velocity. The first is obtained in Section 3.1 by simply arguing as we did in Theorem 2.1. The other, in Section 3.2, achieves an estimate of the perturbed velocity (corresponding to x) in terms of the unperturbed velocity (corresponding to x) via interpolation.

3.1. Direct Bounds. Following closely the arguments in Theorem 2.1, we obtain the next result.

Theorem 3.1. *Fix $W \in C^2(\mathbb{R})$ with $W'' \in L^\infty(\mathbb{R})$, $\mu > 0$, and let $x \in M_b$. There exist numbers $C^w = C^w(x, W)$ and $v^w = v^w(\mu, x)$ for which given any $n, m \in \mathbb{Z}$, the bound*

$$(3.5) \quad \max \left[\left| \frac{\partial}{\partial z} a_n^w(t) \right|, \left| \frac{\partial}{\partial z} b_n^w(t) \right| \right] \leq C^w e^{-\mu(|n-m|-v^w|t|)},$$

holds for all $t \in \mathbb{R}$ and each $z \in \{a_m, b_m\}$. In fact, one may take

$$(3.6) \quad v^w = \left(1 + \sqrt{17 + \frac{4C_2\|W''\|_\infty}{C_1}} \right) C_1 \left(e^{\mu+1} + \frac{1}{\mu} \right)$$

with C_1, C_2 from (3.4).

Proof. Fix $x \in M_b$. Global existence again guarantees that for each $n \in \mathbb{Z}$, the function

$$(3.7) \quad F_n^w(t) = \begin{pmatrix} a_n^w(t) \\ b_n^w(t) \end{pmatrix}$$

satisfies

$$(3.8) \quad F_n^w(t) = F_n^w(0) + \int_0^t \begin{pmatrix} a_n^w(s) (b_{n+1}^w(s) - b_n^w(s)) \\ 2(a_n^w(s)^2 - a_{n-1}^w(s)^2) + R_n(s) \end{pmatrix} ds.$$

Since W is sufficiently smooth, the components of $F_n^w(t)$ are differentiable with respect to each $z \in \{a_m, b_m\}$, and the bound

$$(3.9) \quad \left| \frac{\partial}{\partial z} F_n^w(t) \right| \leq \left| \frac{\partial}{\partial z} F_n^w(0) \right| + \sum_{|e| \leq 1} \int_0^{|t|} D_e^w \left| \frac{\partial}{\partial z} F_{n+e}^w(s) \right| ds,$$

follows as in (2.22) with

$$(3.10) \quad D_e^w = \begin{pmatrix} 2C_1\delta_0(e) & C_1(\delta_0(e) + \delta_1(e)) \\ (4C_1 + C_2\|W''\|_\infty)(\delta_0(e) + \delta_{-1}(e)) & 0 \end{pmatrix}.$$

Following the previous scheme, iteration (with $z = a_m$) yields

$$(3.11) \quad \left| \frac{\partial}{\partial a_m} F_n^w(t) \right| \leq \sum_{k=|n-m|}^{\infty} \frac{(2C_1|t|)^k}{k!} (D^w)^k \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

with

$$(3.12) \quad D^w = \begin{pmatrix} 1 & 1 \\ \alpha & 0 \end{pmatrix} \quad \text{for} \quad \alpha = 4 + \frac{C_2\|W''\|_\infty}{C_1}.$$

Eigenvalues and eigenvectors of D^w are

$$(3.13) \quad \lambda_{\pm} = \frac{1 \pm \sqrt{1+4\alpha}}{2} \quad \text{and} \quad v_{\pm} = \begin{pmatrix} \lambda_{\pm} \\ \alpha \end{pmatrix}$$

and arguing as before, it is now clear that

$$(3.14) \quad \left\| \frac{\partial}{\partial a_m} F_n^w(t) \right\|_{\infty} \leq \frac{2\alpha}{\sqrt{1+4\alpha}} e^{-\mu(|n-m|-v^w|t|)}$$

with

$$(3.15) \quad v^w = 2\lambda_+ C_1 \left(e^{\mu+1} + \frac{1}{\mu} \right).$$

For $z = b_m$, one similarly finds

$$(3.16) \quad \left\| \frac{\partial}{\partial b_m} F_n^w(t) \right\|_{\infty} \leq \frac{2\lambda_+}{\sqrt{1+4\alpha}} e^{-\mu(|n-m|-v^w|t|)}$$

This completes the proof. \square

3.2. Bounds Via Interpolation. The goal of this section is to prove a different bound on the Lieb–Robinson velocity corresponding to an initial condition $x \in M_b$, see Theorem 3.3 below. The novel feature of this estimate is that it is explicit in the unperturbed Lieb–Robinson velocity of x .

Before we state the main result of this section, we first indicate some further estimates on the unperturbed system which will be useful in proving Theorem 3.3. To start with, we prove another Lieb–Robinson type estimate for the Toda system, see Lemma 3.2 below. Afterwards, we introduce a quantity that is better suited for the iteration scheme which is at the heart of proving Theorem 3.3, see $G_\mu(k)$ in (3.37) below. Lastly, we will state and prove Theorem 3.3.

We begin with the following lemma on second order derivatives of the unperturbed system.

Lemma 3.2. *Fix $\mu > 0$ and let $x \in M$. There exists a number $C = C(\mu, x) > 0$ and a function h , depending on μ and x , for which given any $n, k, \ell \in \mathbb{Z}$, the estimate*

$$(3.17) \quad \max \left[\left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} a_n(t) \right|, \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} b_n(t) \right| \right] \leq C e^{-\mu|n-\ell|} e^{-\mu|n-k|} e^{2\mu v|t|} h(t)$$

holds for all $t \in \mathbb{R}$. Here $a_n(t)$ and $b_n(t)$ are the solutions of (2.4) with initial condition $x \in M$, $\frac{\partial}{\partial \tilde{b}_k} = \frac{\partial}{\partial b_{k+1}} - \frac{\partial}{\partial b_k}$, $z \in \{a_\ell, b_\ell\}$, and the number v is as in Theorem 2.1. The function h grows at most exponentially.

Proof. Fix $x \in M$. As in Theorem 2.1, it is clear that the function $F_n(t)$ is well-defined for each $n \in \mathbb{Z}$ and $t \in \mathbb{R}$. In addition, both (2.18) and (2.19) still hold with the choice $z = \tilde{b}_k$. Taking a second derivative, we find that

$$(3.18) \quad \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) = \sum_{|e| \leq 1} \int_0^t \left(\frac{\partial}{\partial z} D_{n,e}(s) \frac{\partial}{\partial \tilde{b}_k} F_{n+e}(s) + D_{n,e}(s) \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_{n+e}(s) \right) ds,$$

for each $z \in \{a_\ell, b_\ell\}$, since $\frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(0) = 0$.

The bound

$$(3.19) \quad \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) \right| \leq \frac{8}{\sqrt{17}} e^{-\mu|n-\ell|} \sum_{|e| \leq 1} \int_0^{|t|} e^{\mu v s} D'_e \left| \frac{\partial}{\partial \tilde{b}_k} F_{n+e}(s) \right| ds + \\ + \|L(0)\|_2 \sum_{|e| \leq 1} \int_0^{|t|} D_e \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_{n+e}(s) \right| ds$$

with

$$(3.20) \quad D'_e = \begin{pmatrix} (e^\mu + 1)\delta_0(e) & \delta_0(e) + \delta_1(e) \\ 4(\delta_0(e) + e^\mu \delta_{-1}(e)) & 0 \end{pmatrix}$$

and D_e as in (2.23), follows using Theorem 2.1 and the argument therein. A further application of Theorem 2.1 implies that

$$(3.21) \quad \left| \frac{\partial}{\partial \tilde{b}_k} F_{n+e}(s) \right| \leq \frac{8}{\sqrt{17}} e^\mu e^{-\mu|n+e-k|} e^{\mu v s} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

and therefore the first term on the right hand side of (3.19) can be estimated by

$$(3.22) \quad \frac{64}{17} e^\mu e^{-\mu|n-\ell|} e^{-\mu|n-k|} \int_0^{|t|} e^{2\mu v s} ds \begin{pmatrix} 2(e^\mu + 1) \\ 4(e^{2\mu} + 1) \end{pmatrix}.$$

Here we have used that

$$(3.23) \quad \sum_{|e| \leq 1} e^{-\mu|n+e-k|} D'_e \begin{pmatrix} 1 \\ 1 \end{pmatrix} \leq e^{-\mu|n-k|} \begin{pmatrix} 2(e^\mu + 1) \\ 4(e^{2\mu} + 1) \end{pmatrix}.$$

Let us introduce the notation

$$(3.24) \quad C_\mu = \frac{64}{17} e^\mu \quad \text{and} \quad y = \begin{pmatrix} 2(e^\mu + 1) \\ 4(e^{2\mu} + 1) \end{pmatrix}.$$

We have proven that

$$(3.25) \quad \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) \right| \leq C_\mu e^{-\mu|n-\ell|} e^{-\mu|n-k|} \int_0^{|t|} e^{2\mu v s} ds \cdot y + \|L(0)\|_2 \int_0^{|t|} \sum_{|e| \leq 1} D_e \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_{n+e}(s) \right| ds$$

Iteration now yields

$$(3.26) \quad \begin{aligned} \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) \right| &\leq C_\mu \sum_{j=0}^{\infty} \|L(0)\|_2^j \int_0^{|t|} \int_0^{t_1} \dots \int_0^{t_j} e^{2\mu v t_{j+1}} dt_{j+1} \dots dt_1 \times \\ &\quad \times \sum_{|e_1| \leq 1} \dots \sum_{|e_j| \leq 1} e^{-\mu|n+e_1+\dots+e_j-\ell|} e^{-\mu|n+e_1+\dots+e_j-k|} D_{e_1} D_{e_2} \dots D_{e_j} y \end{aligned}$$

and the convergence is guaranteed as before.

A short calculation shows that

$$(3.27) \quad \int_0^{|t|} \int_0^{t_1} \dots \int_0^{t_j} e^{2\mu v t_{j+1}} dt_{j+1} \dots dt_1 = \frac{1}{(2\mu v)^{j+1}} \sum_{m=j+1}^{\infty} \frac{(2\mu v |t|)^m}{m!}$$

whereas the bound $e^{-\mu|z+e|} \leq e^{\mu|e|} e^{-\mu|z|}$ immediately implies that

$$(3.28) \quad e^{-\mu|n+e_1+e_2+\dots+e_j-\ell|} \leq e^{\mu|e_1|} \dots e^{\mu|e_j|} e^{-\mu|n-\ell|}.$$

This proves that

$$(3.29) \quad \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) \right| \leq \frac{C_\mu}{2\mu v} e^{-\mu|n-\ell|} e^{-\mu|n-k|} \sum_{j=0}^{\infty} \left(\frac{\|L(0)\|_2}{2\mu v} \right)^j \sum_{m=j+1}^{\infty} \frac{(2\mu v |t|)^m}{m!} (\tilde{D})^j y$$

where

$$(3.30) \quad \tilde{D} = \sum_{|e| \leq 1} e^{2\mu|e|} D_e = \begin{pmatrix} 2 & e^{2\mu} + 1 \\ 4(e^{2\mu} + 1) & 0 \end{pmatrix}.$$

The eigenvalues of \tilde{D} are $\lambda_{\pm} = 1 \pm \sqrt{1 + 4(e^{2\mu} + 1)^2}$ and a convenient choice of eigenvectors are

$$(3.31) \quad v_{\pm} = \begin{pmatrix} \lambda_{\pm} \\ 4(e^{2\mu} + 1) \end{pmatrix}.$$

In terms of these, it is clear that

$$(3.32) \quad y = y_1 v_+ + y_2 v_- \quad \text{with} \quad y_1 = \frac{2(e^{\mu} + 1) - \lambda_-}{\lambda_+ - \lambda_-} \quad \text{and} \quad y_2 = \frac{\lambda_+ - 2(e^{\mu} + 1)}{\lambda_+ - \lambda_-}.$$

Inserting this into (3.29), we find that

$$(3.33) \quad \begin{aligned} \left\| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) \right\|_{\infty} &\leq \frac{C_\mu}{2\mu v} e^{-\mu|n-\ell|} e^{-\mu|n-k|} \sum_{j=0}^{\infty} \left(\frac{\|L(0)\|_2}{2\mu v} \right)^j \sum_{m=j+1}^{\infty} \frac{(2\mu v |t|)^m}{m!} \|y_1 \lambda_+^j v_+ + y_2 \lambda_-^j v_-\|_{\infty} \\ &\leq C e^{-\mu|n-\ell|} e^{-\mu|n-k|} \sum_{j=0}^{\infty} \left(\frac{\|L(0)\|_2 \lambda_+}{2\mu v} \right)^j \sum_{m=j+1}^{\infty} \frac{(2\mu v |t|)^m}{m!} \end{aligned}$$

It is easy to see that for any $A, B \in \mathbb{R}$,

$$(3.34) \quad \sum_{j=0}^{\infty} A^j \sum_{m=j+1}^{\infty} \frac{(B|t|)^m}{m!} = \sum_{m=1}^{\infty} \frac{(B|t|)^m}{m!} \sum_{j=0}^{m-1} A^j = \begin{cases} \frac{1}{A-1} (e^{(A-1)B|t|} - 1) e^{B|t|}, & A \neq 1, \\ B|t| e^{B|t|}, & A = 1. \end{cases}$$

This proves Lemma 3.2. One may choose

$$(3.35) \quad C = \frac{C_\mu}{2\mu v} (|y_1| \|v_+\|_{\infty} + |y_2| \|v_-\|_{\infty})$$

and by setting $\beta = \frac{\|L(0)\|_2 \lambda_+}{2\mu v} - 1$, we see that

$$(3.36) \quad h(t) = \begin{cases} \frac{1}{\beta} (e^{2\mu v \beta |t|} - 1) & \text{if } \beta \neq 0, \\ 2\mu v |t| & \text{if } \beta = 0. \end{cases}$$

Note that β depends only on μ , $\lim_{t \rightarrow 0} h(t) = 0$, and moreover, if $\beta < 0$, then $h(t) \leq |\beta|^{-1}$. For any μ , however, $h(t) \leq h e^{r|t|}$, where $h = h(\mu, x)$ and $r = r(\mu, x) \geq 0$. \square

Our proof of Theorem 3.3 will again use an iteration scheme. Due to interpolation, it will be necessary to sum certain terms over all integers. To this end, we make the following observation. For each $\mu > 0$, the function

$$(3.37) \quad G_\mu(k) = \frac{e^{-\mu|k|}}{(1 + |k|)^2}$$

satisfies the estimate

$$(3.38) \quad \sum_{l \in \mathbb{Z}} G_\mu(j-l) G_\mu(l-k) \leq \gamma G_\mu(j-k)$$

with $\gamma = 4 \sum_{k \in \mathbb{Z}} (1 + |k|)^{-2}$. In fact, setting $f(k) = (1 + |k|)^2$, it is clear that $f(j+k) \leq 2(f(j) + f(k))$, and thus

$$(3.39) \quad \begin{aligned} G_\mu(j-k)^{-1} G_\mu(j-l) G_\mu(l-k) &= f(j-k)^{-1} f(j-l)^{-1} f(l-k)^{-1} e^{\mu(|j-k| - |j-l| - |l-k|)} \\ &\leq 2(f(j-l)^{-1} + f(l-k)^{-1}). \end{aligned}$$

The claim follows after summing over l . Clearly the choice of power 2 in the denominator of G_μ is merely for convenience; a factor $(1 + |k|)^{-1-\delta}$ for any $\delta > 0$ would suffice. Note, however, that there is no bound of the type in (3.38) for the exponential function with no inverse polynomial weight.

As a final comment, it is clear that for any $\mu > 0$ and $\varepsilon > 0$, the bound

$$(3.40) \quad e^{-(\mu+\varepsilon)|x|} \leq C_\varepsilon G_\mu(|x|)$$

holds for all $x \in \mathbb{R}$. The number $C_\varepsilon = \sup_{x \in \mathbb{R}} (1 + |x|)^2 e^{-\varepsilon|x|} < \infty$.

We can now state the second result of this section.

Theorem 3.3. *Fix $W \in C^2(\mathbb{R})$ with $W', W'' \in L^\infty(\mathbb{R})$, take $\mu > 0$, and let $x \in M_b$. For any $\varepsilon > 0$, there are positive numbers $C = C(\varepsilon)$, $D = D(\varepsilon, \mu, x, W)$, and $\delta = \delta(\varepsilon, \mu, x, W)$ such that*

$$(3.41) \quad \max \left[\left| \frac{\partial}{\partial z} a_n^w(t) \right|, \left| \frac{\partial}{\partial z} b_n^w(t) \right| \right] \leq C G_\mu(|n-m|) e^{(\mu+\varepsilon)v|t|} \left[1 + D \left(e^{\delta|t|} - 1 \right) \right]$$

holds for all $t \in \mathbb{R}$ and $n, m \in \mathbb{Z}$. Here $v = v(x, \mu + \varepsilon)$ is as in Theorem 2.1 and $z \in \{a_m, b_m\}$.

In words, this result shows that for each $\mu > 0$ the perturbed Lieb–Robinson velocity $v^w(x, \mu)$, corresponding to $x \in M_b$, satisfies, for each $\varepsilon > 0$,

$$(3.42) \quad v^w(x, \mu) \leq \left(1 + \frac{\varepsilon}{\mu} \right) v(x, \mu + \varepsilon) + \frac{\delta}{\mu}$$

and, as is shown in the proof below, the dependence of δ on W can be made explicit

$$(3.43) \quad \delta = \delta_1 \|W'\|_\infty + \delta_2 \|W''\|_\infty + \delta_3$$

for some $\delta_i = \delta_i(\varepsilon, \mu, x)$ and $i = 1, 2, 3$.

Proof. Fix W , $x \in M_b$, and $t \in \mathbb{R}$. We begin by interpolating between the Toda and perturbed dynamics. As in Section 2.1, denote by

$$(3.44) \quad \alpha_t(A) = A \circ \Phi_t \quad \text{and} \quad \alpha_t^w(A) = A \circ \Phi_t^w \quad \text{for any } A \in \mathcal{A},$$

the Toda and perturbed dynamics respectively. For any $A \in \mathcal{A}$, the observable-valued equation

$$(3.45) \quad \alpha_t^w(A) - \alpha_t(A) = \int_0^t \frac{d}{ds} \alpha_s^w(\alpha_{t-s}(A)) ds$$

is clear. Due to the form of the perturbation,

$$(3.46) \quad \begin{aligned} \frac{d}{ds} \alpha_s^w(\alpha_{t-s}(A)) &= \alpha_s^w(\{\alpha_{t-s}(A), H^w\}) - \alpha_s^w(\alpha_{t-s}(\{A, H\})) \\ &= \sum_{k \in \mathbb{Z}} \alpha_s^w(\{\alpha_{t-s}(A), W_k\}), \end{aligned}$$

and moreover, one finds that

$$(3.47) \quad \{\alpha_{t-s}(A), W_k\} = -\frac{1}{2} W'_k \cdot \frac{\partial}{\partial b_k} \alpha_{t-s}(A).$$

Combining these expressions, we have shown that

$$(3.48) \quad \alpha_t^w(A) = \alpha_t(A) - \frac{1}{2} \sum_{k \in \mathbb{Z}} \int_0^t \alpha_s^w \left(W'_k \cdot \frac{\partial}{\partial b_k} \alpha_{t-s}(A) \right) ds$$

as a formal expansion. For specific choices of observables, the expression above can be estimated.

Fix $n \in \mathbb{Z}$ and denote by A_n and B_n the observables introduced in Example 2.2. Consider the observable-valued

$$(3.49) \quad \tilde{F}_n^w(t) = \begin{pmatrix} \alpha_t^w(A_n) \\ \alpha_t^w(B_n) \end{pmatrix} \quad \text{and} \quad \tilde{F}_n(t) = \begin{pmatrix} \alpha_t(A_n) \\ \alpha_t(B_n) \end{pmatrix},$$

which for any given $x \in M$ satisfies

$$(3.50) \quad [\tilde{F}_n^w(t)](x) = \begin{pmatrix} a_n^w(t) \\ b_n^w(t) \end{pmatrix} = F_n^w(t) \quad \text{and} \quad [\tilde{F}_n(t)](x) = \begin{pmatrix} a_n(t) \\ b_n(t) \end{pmatrix} = F_n(t),$$

in terms of our previous notation. Our starting point is the equation

$$(3.51) \quad \tilde{F}_n^w(t) = \tilde{F}_n(t) - \frac{1}{2} \sum_{k \in \mathbb{Z}} \int_0^t \alpha_s^w(W'_k) \cdot \begin{pmatrix} \alpha_s^w \left(\frac{\partial}{\partial b_k} \alpha_{t-s}(A_n) \right) \\ \alpha_s^w \left(\frac{\partial}{\partial b_k} \alpha_{t-s}(B_n) \right) \end{pmatrix} ds.$$

Since W is sufficiently smooth, the formula

$$(3.52) \quad \frac{\partial}{\partial z} \tilde{F}_n^w(t) = \frac{\partial}{\partial z} \tilde{F}_n(t) - \sum_{k \in \mathbb{Z}} \int_0^t D_k(s; n) \frac{\partial}{\partial z} \tilde{F}_k^w(s) ds - \frac{1}{2} \sum_{k, \ell \in \mathbb{Z}} \int_0^t \alpha_s^w(W'_k) \cdot D_{k, \ell}(s; n) \frac{\partial}{\partial z} \tilde{F}_\ell^w(s) ds$$

where

$$(3.53) \quad D_k(s; n) = \begin{pmatrix} \alpha_s^w(W'_k \cdot A_k^{-1}) \cdot \alpha_s^w \left(\frac{\partial}{\partial b_k} \alpha_{t-s}(A_n) \right) & 0 \\ \alpha_s^w(W'_k \cdot A_k^{-1}) \cdot \alpha_s^w \left(\frac{\partial}{\partial b_k} \alpha_{t-s}(B_n) \right) & 0 \end{pmatrix}$$

and

$$(3.54) \quad D_{k, \ell}(s; n) = \begin{pmatrix} \alpha_s^w \left(\frac{\partial^2}{\partial a_\ell \partial b_k} \alpha_{t-s}(A_n) \right) & \alpha_s^w \left(\frac{\partial^2}{\partial b_\ell \partial b_k} \alpha_{t-s}(A_n) \right) \\ \alpha_s^w \left(\frac{\partial^2}{\partial a_\ell \partial b_k} \alpha_{t-s}(B_n) \right) & \alpha_s^w \left(\frac{\partial^2}{\partial b_\ell \partial b_k} \alpha_{t-s}(B_n) \right) \end{pmatrix}$$

readily follows.

For any $x \in M$, the estimate

$$(3.55) \quad \begin{aligned} \left\| \left[\frac{\partial}{\partial z} \tilde{F}_n^w(t) \right] (x) \right\|_\infty &\leq \left\| \left[\frac{\partial}{\partial z} \tilde{F}_n(t) \right] (x) \right\|_\infty + \sum_{k \in \mathbb{Z}} \int_0^{|t|} \left\| \left[D_k(s; n) \frac{\partial}{\partial z} \tilde{F}_k^w(s) \right] (x) \right\|_\infty ds \\ &\quad + \frac{\|W'\|_\infty}{2} \sum_{k, \ell \in \mathbb{Z}} \int_0^{|t|} \left\| \left[D_{k, \ell}(s; n) \frac{\partial}{\partial z} \tilde{F}_\ell^w(s) \right] (x) \right\|_\infty ds \end{aligned}$$

is clear.

Using Theorem 2.1, the first term can be bounded by

$$(3.56) \quad \begin{aligned} \left\| \left[\frac{\partial}{\partial z} \tilde{F}_n(t) \right] (x) \right\|_\infty &= \left\| \frac{\partial}{\partial z} F_n(t) \right\|_\infty \\ &\leq \frac{8}{\sqrt{17}} e^{-(\mu+\varepsilon)(|n-m|-v|t|)} \\ &\leq \frac{8}{\sqrt{17}} C_\varepsilon G_\mu (|n-m|) e^{(\mu+\varepsilon)v|t|} \end{aligned}$$

where we have set $v = v(x, \mu + \varepsilon)$ and used (3.40).

For each $x \in M_b$, the matrix appearing in the second term satisfies

$$(3.57) \quad \begin{aligned} \|[D_k(s; n)](x)\|_\infty &\leq C_2 \|W''\|_\infty \frac{8}{\sqrt{17}} e^{\mu+\varepsilon} e^{-(\mu+\varepsilon)(|n-k|-v(s)(|t|-s))} \\ &\leq C_2 \|W''\|_\infty \frac{8}{\sqrt{17}} e^{\mu+\varepsilon} C_\varepsilon G_\mu(|n-k|) e^{(\mu+\varepsilon)v(s)(|t|-s)} \end{aligned}$$

where we have denoted by

$$(3.58) \quad v(s) = v(\Phi_s^w(x), \mu + \varepsilon) = (1 + \sqrt{17}) \|L(\Phi_s^w(x))\|_2 (e^{\mu+\varepsilon+1} + (\mu + \varepsilon)^{-1})$$

the unperturbed Lieb–Robinson corresponding to the initial condition $\Phi_s^w(x)$. Since $x \in M_b$,

$$(3.59) \quad \|L(\Phi_s^w(x))\|_2 \leq 2 \|a^w(s)\|_\infty + \|b^w(s)\|_\infty \leq 3C_1,$$

and so the quantity in (3.58) can be estimated independent of s . Clearly,

$$(3.60) \quad v = v(0) \leq \sup_{s \in \mathbb{R}} v(s) := v^*.$$

The matrix in the third term can be dominated using Lemma 3.2. In fact, the bound

$$(3.61) \quad \|[D_{k,\ell}(s; n)](x)\|_\infty \leq 2C e^{-(\mu+\varepsilon)|n-\ell|} e^{-(\mu+\varepsilon)|n-k|} e^{2(\mu+\varepsilon)v(s)(|t|-s)} h(|t|-s)$$

follows immediately. Some comments are in order. First, the prefactor C from Lemma 3.2, appearing above, seems to depend on s through the velocity $v(s)$; see (3.35). If, however, one repeats the argument of Lemma 3.2 and replaces the $v(s)$ in (3.27) with v^* above, then the new prefactor is independent of s . Next, as discussed at the end of Lemma 3.2, the function h grows at most exponentially. In fact, in this particular application, it is clear that

$$(3.62) \quad h(|t|-s) \leq h e^{2(\mu+\varepsilon)\gamma v^*(|t|-s)}$$

with numbers $h = h(x, \mu + \varepsilon)$ and $\gamma = \gamma(x, \mu + \varepsilon)$, each independent of s . This proves that

$$(3.63) \quad \|[D_{k,\ell}(s; n)](x)\|_\infty \leq \tilde{C} G_\mu(|n-\ell|) G_\mu(|n-k|) e^{2(\gamma+1)(\mu+\varepsilon)v^*(|t|-s)}.$$

Putting everything together and suppressing the x -dependence, we have found that

$$(3.64) \quad \begin{aligned} \left\| \frac{\partial}{\partial z} F_n^w(t) \right\|_\infty &\leq B_1 G_\mu(|n-m|) e^{(\mu+\varepsilon)v|t|} + B_2 \sum_{k \in \mathbb{Z}} G_\mu(|n-k|) \int_0^{|t|} e^{(\mu+\varepsilon)v^*(|t|-s)} \left\| \frac{\partial}{\partial z} F_k^w(s) \right\|_\infty ds \\ &\quad + B_3 \sum_{k, \ell \in \mathbb{Z}} G_\mu(|n-\ell|) G_\mu(|n-k|) \int_0^{|t|} e^{2(\gamma+1)(\mu+\varepsilon)v^*(|t|-s)} \left\| \frac{\partial}{\partial z} F_\ell^w(s) \right\|_\infty ds \\ &\leq B_1 G_\mu(|n-m|) e^{(\mu+\varepsilon)v|t|} + B \sum_{k \in \mathbb{Z}} G_\mu(|n-k|) \int_0^{|t|} e^{2(\gamma+1)(\mu+\varepsilon)v^*(|t|-s)} \left\| \frac{\partial}{\partial z} F_k^w(s) \right\|_\infty ds \end{aligned}$$

Upon iteration, we find that

$$\begin{aligned} \left\| \frac{\partial}{\partial z} F_n^w(t) \right\|_\infty &\leq B_1 G_\mu(|n-m|) e^{a|t|} + B_1 \sum_{j=1}^{\infty} B^j \sum_{k_1 \in \mathbb{Z}} \cdots \sum_{k_j \in \mathbb{Z}} G_\mu(|n-k_1|) \cdots G_\mu(|k_j-m|) \\ &\quad \times \int_0^{|t|} \int_0^{s_1} \cdots \int_0^{s_{j-1}} e^{b(|t|-s_1)} e^{b(s_1-s_2)} \cdots e^{b(s_{j-1}-s_j)} e^{a s_j} ds_j \cdots ds_1 \end{aligned}$$

where we have set

$$(3.65) \quad a = (\mu + \varepsilon)v \quad \text{and} \quad b = 2(\gamma + 1)(\mu + \varepsilon)v^*.$$

From (3.38), it is clear that

$$(3.66) \quad \sum_{k_1 \in \mathbb{Z}} \cdots \sum_{k_j \in \mathbb{Z}} G_\mu(|n-k_1|) \cdots G_\mu(|k_j-m|) \leq \gamma^j G_\mu(|n-m|)$$

and the iterated integral can also be calculated:

$$\begin{aligned}
 \int_0^{|t|} \int_0^{s_1} \dots \int_0^{s_{j-1}} e^{b(|t|-s_1)} e^{b(s_1-s_2)} \dots e^{b(s_{j-1}-s_j)} e^{as_j} ds_j \dots ds_1 &= e^{b|t|} \int_0^{|t|} \int_0^{s_1} \dots \int_0^{s_{j-1}} e^{(a-b)s_j} ds_j \dots ds_1 \\
 (3.67) \qquad \qquad \qquad &= \frac{e^{b|t|}}{(a-b)^j} \sum_{k=j}^{\infty} \frac{((a-b)|t|)^k}{k!}
 \end{aligned}$$

This shows that

$$(3.68) \qquad \left\| \frac{\partial}{\partial z} F_n^w(t) \right\|_{\infty} \leq B_1 G_{\mu}(|n-m|) \left(e^{a|t|} + e^{b|t|} \sum_{j=1}^{\infty} \left(\frac{B\gamma}{a-b} \right)^j \sum_{k=j}^{\infty} \frac{((a-b)|t|)^k}{k!} \right)$$

Since $v_{\mu+\varepsilon} \leq v^*$, it is clear that $a-b < 0$. In this case,

$$(3.69) \qquad \sum_{j=1}^{\infty} \left(\frac{B\gamma}{a-b} \right)^j \sum_{k=j}^{\infty} \frac{((a-b)|t|)^k}{k!} = \frac{B\gamma}{B\gamma + b - a} \left[e^{B\gamma|t|} - e^{(a-b)|t|} \right]$$

follows from (3.34) and so

$$(3.70) \qquad \left\| \frac{\partial}{\partial z} F_n^w(t) \right\|_{\infty} \leq B_1 G_{\mu}(|n-m|) e^{a|t|} \left(1 + \frac{B\gamma}{B\gamma + b - a} \left[e^{(B\gamma+b-a)|t|} - 1 \right] \right).$$

This proves (3.41) and completes the proof. \square

3.3. Comments on these results. It is clear that results analogous to Theorem 3.1 and Theorem 3.3 hold for more general, finite range potentials W . We stated the results as above for simplicity of presentation.

It is also clear that analogues of Theorem 2.3, for more general observables, hold as direct corollaries of Theorem 3.1 and Theorem 3.3. Since the statements are clear, we do not rewrite them for the sake of brevity.

4. BOUNDS FOR THE HIERARCHY

In this section, we will demonstrate that the results from Section 2 also apply to the Toda hierarchy. We begin, in Section 4.1, by introducing the Toda hierarchy using the recursive approach from [8]. We will use essentially the same notation as [40] and refer the interested reader to this text for further information. In Section 4.2, we state the Lieb–Robinson bound valid for the hierarchy, see Theorem 4.4. To prove this result, we first establish the crucial solution estimate, an analogue of Theorem 2.1, which shows that a solution of the hierarchy at site n has weak dependence on the initial condition at site m if $|n-m| \gg 1$. The bound is explicit in terms of a quantity v_r which, in particular, depends on the initial condition through the infinity norm of a specific matrix, see (4.11). To make our bound even more concrete, we estimate the norm of this matrix in Lemma 4.3 below. Theorem 4.4 follows as in Section 2.

We note that our results also immediately apply to the Kac–van Moerbeke hierarchy. This hierarchy can be viewed as a special case of the Toda hierarchy which is obtained by setting $b_n \equiv 0$ in the even order Toda equations. As is discussed, e.g. in [27], this gives precisely the equations of the Kac–van Moerbeke hierarchy.

4.1. The Toda Hierarchy. We introduce the Toda Hierarchy as follows. Fix $r \in \mathbb{N}_0 = \mathbb{N} \cup \{0\}$. Take $c_0 = 1$ and choose constants $c_j \in \mathbb{R}$ for $1 \leq j \leq r$. For each $x \in M$, define sequences $g(x, r)$ and $h(x, r)$ componentwise by setting

$$(4.1) \qquad g_n(x, r) = \sum_{j=0}^r c_{r-j} \tilde{g}_n^{(j+1)}(x) \quad \text{with} \quad \tilde{g}_n^{(j)}(x) = \langle \delta_n, L(x)^j \delta_n \rangle,$$

and

$$(4.2) \qquad h_n(x, r) = \sum_{j=0}^r c_{r-j} \tilde{h}_n^{(j+1)}(x) \quad \text{with} \quad \tilde{h}_n^{(j)}(x) = 2a_n \langle \delta_{n+1}, L(x)^j \delta_n \rangle,$$

with $L(x)$ as in (2.8). We now define a system of equations for the components of an unknown sequence $x(t, r) = \{(a_n(t, r), b_n(t, r))\}$:

$$(4.3) \qquad \dot{a}_n(t, r) = a_n(t, r) (g_{n+1}(t, r) - g_n(t, r))$$

and

$$(4.4) \quad \dot{b}_n(t, r) = (h_n(t, r) - h_{n-1}(t, r))$$

with initial condition $x(0) = \{(a_n(0, r), b_n(0, r))\} = x \in M$. As in the previous section, we have, for example, denoted by $g_n(t, r) = g_n(x(t, r), r)$ to simplify notation. The system (4.3) and (4.4) is known to have global solutions (see e.g. [40], Theorem 12.6) for initial conditions in M . Varying $r \in \mathbb{N}_0$ describes the Toda hierarchy. Let us denote by $\Phi_t^{(r)}$ the flow corresponding to the Toda hierarchy, i.e. $\Phi_t^{(r)} : M \rightarrow M$ satisfies $\Phi_t^{(r)}(x) = \{(a_n(t, r), b_n(t, r))\}$, the solution of (4.3) and (4.4) above, with $\Phi_0^{(r)}(x) = x$.

The simplest example corresponds to $r = 0$. In this case, the system becomes

$$(4.5) \quad \dot{a}_n(t, 0) = a_n(t, 0)(b_{n+1}(t, 0) - b_n(t, 0)) \quad \text{and} \quad \dot{b}_n(t, 0) = 2(a_n(t, 0)^2 - a_{n-1}(t, 0)^2),$$

which is, of course, the Toda system (2.4).

4.2. The Lieb–Robinson Bound. The above choice of sequences g and h guarantee that the Lax-formalism of the Toda Lattice, see (2.8)–(2.12), still holds for the hierarchy. In fact, set

$$(4.6) \quad P(x, r) = \sum_{j=0}^r c_{r-j} \tilde{P}^{(j+1)}(x) \quad \text{with} \quad \tilde{P}^{(j)}(x) = [L(x)^j]_+ - [L(x)^j]_-,$$

where $[A]_{\pm}$ denote the upper and lower triangular parts of an operator with respect to the standard basis $\delta_m(n) = \delta_{m,n}$ (with $\delta_{m,n}$ the usual Kronecker delta). It is known, see e.g. [40], that the Toda hierarchy is equivalent to the Lax equation

$$(4.7) \quad \frac{d}{dt} L(t) = [P(t, r), L(t)],$$

where we have set $L(t) = L(\Phi_t^{(r)}(x))$ and $P(t, r) = P(\Phi_t^{(r)}(x), r)$ again to ease notation. It is easy to see that the operator $P(t, r)$, which is of order $2r + 2$, is skew-adjoint and differentiable. Like before then, there exists a unique unitary propagator $U^{(r)}(t, s)$ for $P(t, r)$. It follows from the Lax equation that

$$(4.8) \quad L(t) = U^{(r)}(t, s) L(s) U^{(r)}(t, s)^{-1}$$

implying again an a-priori estimate

$$(4.9) \quad \max(\|a(t, r)\|_{\infty}, \|b(t, r)\|_{\infty}) \leq \|L(t)\|_2 = \|L(0)\|_2.$$

For later use we record the following structure.

Lemma 4.1 ([23]). *For each $x \in M$ and any integer $j \geq 1$, the sequences $\tilde{g}^{(j)}(x)$ and $\tilde{h}^{(j)}(x)$, as defined in (4.1) and (4.2), have components that are homogeneous. In fact, they are sums of monomials of the components of x with degree j and $j + 1$, respectively, which have the form:*

$$\tilde{g}_n^{(j)}(x) = \begin{cases} \left(\prod_{\ell=0}^{k-1} a_{n+\ell}^2 \right) b_{n+k} + R(n+k-1, n-k+1) + \\ + \left(\prod_{\ell=1}^k a_{n-\ell}^2 \right) \left(b_{n-k} + 2 \sum_{\ell=0}^{k-1} b_{n-\ell} \right), & j = 2k+1, \\ \left(\prod_{\ell=0}^{k-2} a_{n+\ell}^2 \right) \left(a_{n+k-1}^2 + b_{n+k-1}^2 + 2b_{n+k-1} \sum_{\ell=0}^{k-2} b_{n+\ell} \right) + \\ + R(n+k-2, n-k+1) + \prod_{\ell=1}^k a_{n-\ell}^2, & j = 2k, \end{cases}$$

and

$$\tilde{h}_n^{(j)}(x) = \begin{cases} 2 \left(\prod_{\ell=0}^{k-1} a_{n+\ell}^2 \right) \left(a_{n+k}^2 + b_{n+k}^2 + 2b_{n+k} \sum_{\ell=0}^{k-1} b_{n+\ell} \right) + \\ + R(n+k-1, n-k+1) + 2 \prod_{\ell=0}^k a_{n-\ell}^2, & j = 2k+1, \\ 2 \left(\prod_{\ell=0}^{k-1} a_{n+\ell}^2 \right) b_{n+k} + R(n+k-1, n-k+2) + \\ + 2 \prod_{\ell=0}^{k-1} a_{n-\ell}^2 \left(b_{n+1} + b_{n-k+1} + 2 \sum_{\ell=0}^{k-2} b_{n-\ell} \right), & j = 2k, \end{cases}$$

for $j > 1$. Here $R(n, m)$ denotes terms which involve only a_ℓ and b_ℓ with $m \leq \ell \leq n$ and we set $R(n, m) = 0$ if $n < m$.

The result below is an analogue of Theorem 2.1 for the Toda hierarchy.

Theorem 4.2. Fix $x \in M$, $\mu > 0$ and $r \in \mathbb{N}_0$. For any $n, m \in \mathbb{Z}$, the bound

$$(4.10) \quad \max \left[\left| \frac{\partial}{\partial z} a_n(t, r) \right|, \left| \frac{\partial}{\partial z} b_n(t, r) \right| \right] \leq e^{-\mu(\lceil |n-m|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil - v_r |t|)},$$

holds for all $t \in \mathbb{R}$, where $z \in \{a_m, b_m\}$,

$$(4.11) \quad v_r = v_r(x, \mu) = \|D(r)\|_\infty \|L(0)\|_2 \left(e^{\mu+1} + \frac{1}{\mu} \right),$$

and $D(r)$ depends on x , r , and the numbers c_1, \dots, c_r .

Proof. Without loss of generality assume $t \geq 0$. For each $n \in \mathbb{Z}$, define the function $F_n : \mathbb{R} \rightarrow \mathbb{R}^2$ by

$$(4.12) \quad F_n(t, r) = \begin{pmatrix} a_n(t, r) \\ b_n(t, r) \end{pmatrix}.$$

Since r will be fixed for the remainder of the argument, we will drop it from our notation. Using the equations of motion (4.3) and (4.4), it is clear that

$$(4.13) \quad \begin{aligned} F_n(t) &= F_n(0) + \int_0^t \begin{pmatrix} a_n(s)(g_{n+1}(s) - g_n(s)) \\ h_n(s) - h_{n-1}(s) \end{pmatrix} ds \\ &= F_n(0) + \sum_{j=0}^r c_{r-j} \int_0^t \begin{pmatrix} a_n(s)(\tilde{g}_{n+1}^{(j+1)}(s) - \tilde{g}_n^{(j+1)}(s)) \\ \tilde{h}_n^{(j+1)}(s) - \tilde{h}_{n-1}^{(j+1)}(s) \end{pmatrix} ds \end{aligned}$$

Observe that by Lemma 4.1, for each $0 \leq j \leq r$, the quantities $a_n(s)(\tilde{g}_{n+1}^{(j+1)}(s) - \tilde{g}_n^{(j+1)}(s))$ and $\tilde{h}_n^{(j+1)}(s) - \tilde{h}_{n-1}^{(j+1)}(s)$ are homogeneous polynomials (of degree $j+2$) in the variables $a_{n+e}(s)$ and $b_{n+e}(s)$ for $|e| \leq \lfloor \frac{j}{2} \rfloor + 1$.

Differentiating with respect to $z \in \{a_m, b_m\}$ we get

$$(4.14) \quad \frac{\partial}{\partial z} F_n(t) = \frac{\partial}{\partial z} F_n(0) + \sum_{j=0}^r c_{r-j} \sum_{|e| \leq \lfloor \frac{j}{2} \rfloor + 1} \int_0^t D_{n,e}^{(j)}(s) \frac{\partial}{\partial z} F_{n+e}(s) ds,$$

where the entries of $D_{n,e}^{(j)}(s)$ are homogeneous polynomials of degree $j+1$. In particular, the following estimate, analogous to the bound in (2.22), holds

$$(4.15) \quad |D_{n,e}^{(j)}(s)| \leq \|L(0)\|_2^{j+1} D_e^{(j)},$$

an explicit formula for $D_e^{(j)}$ appears in (4.27) below. As a result, it is clear that

$$(4.16) \quad \left| \frac{\partial}{\partial z} F_n(t) \right| \leq \left| \frac{\partial}{\partial z} F_n(0) \right| + \sum_{j=0}^r |c_{r-j}| \|L(0)\|_2^{j+1} \sum_{|e| \leq \lfloor \frac{j}{2} \rfloor + 1} \int_0^t D_e^{(j)} \left| \frac{\partial}{\partial z} F_{n+e}(s) \right| ds.$$

Upon iteration of the above inequality (in the case that $z = a_m$), we find that

$$(4.17) \quad \begin{aligned} \left| \frac{\partial}{\partial a_m} F_n(t) \right| &\leq \sum_{k=0}^{\infty} \frac{t^k}{k!} \sum_{j_1=0}^r |c_{r-j_1}| \|L(0)\|_2^{j_1+1} \dots \sum_{j_k=0}^r |c_{r-j_k}| \|L(0)\|_2^{j_k+1} \times \\ &\quad \times \sum_{|e_1| \leq \lfloor \frac{j_1}{2} \rfloor + 1} \dots \sum_{|e_k| \leq \lfloor \frac{j_k}{2} \rfloor + 1} \delta_{m+e_1+\dots+e_k}(n) D_{e_1}^{(j_1)} \dots D_{e_k}^{(j_k)} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ &\leq \sum_{k=\lceil |n-m|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil}^{\infty} \frac{(\|L(0)\|_2 t)^k}{k!} \sum_{j_1=0}^r |c_{r-j_1}| \|L(0)\|_2^{j_1} \dots \sum_{j_k=0}^r |c_{r-j_k}| \|L(0)\|_2^{j_k} D^{(j_1)} \dots D^{(j_k)} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ &\leq \sum_{k=\lceil |n-m|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil}^{\infty} \frac{(\|L(0)\|_2 t)^k}{k!} D(r)^k \begin{pmatrix} 1 \\ 0 \end{pmatrix} \end{aligned}$$

where we have set

$$(4.18) \quad D^{(j)} = \sum_{|e| \leq \lfloor \frac{j}{2} \rfloor + 1} D_e^{(j)} \quad \text{and} \quad D(r) = \sum_{j=0}^r |c_{r-j}| \|L(0)\|_2^j D^{(j)}.$$

A similar estimate holds for the case $z = b_m$.

Taking the infinity norm one obtains that for any $\mu > 0$

$$(4.19) \quad \left\| \frac{\partial}{\partial z} F_n(t) \right\|_\infty \leq \sum_{k=\lceil |n-m|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil}^{\infty} \frac{(\|L(0)\|_2 t)^k}{k!} \|D(r)\|_\infty^k \\ \leq e^{-\mu(\lceil |n-m|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil - v_r |t|)},$$

where we have set $v_r = \|D(r)\|_\infty \|L(0)\|_2 \left(e^{\mu+1} + \frac{1}{\mu}\right)$. \square

We now provide a rough, but explicit, estimate on the velocity corresponding to the Toda hierarchy.

Lemma 4.3. *Fix $x \in M$, $\mu > 0$ and $r \in \mathbb{N}_0$. The velocity corresponding to the Toda hierarchy, see (4.10) and (4.11) satisfies*

$$(4.20) \quad v_r \leq 8 \left(e^{\mu+1} + \frac{1}{\mu} \right) \sum_{j=0}^r |c_{r-j}| \|L(0)\|_2^{j+1} (j+2) 3^j$$

Proof. It is clear from (4.11) that we need only estimate the quantity $\|D(r)\|_\infty$, with $D(r)$ as defined in (4.18). To see this, first recall that for each $1 \leq j \leq r$, the quantities $a_n(s)(\tilde{g}_{n+1}^{(j+1)}(s) - \tilde{g}_n^{(j+1)}(s))$ and $\tilde{h}_n^{(j+1)}(s) - \tilde{h}_{n-1}^{(j+1)}(s)$ are homogeneous polynomials (of degree $j+2$) in the variables $a_{n+e}(s)$ and $b_{n+e}(s)$ for $|e| \leq \lfloor \frac{j}{2} \rfloor + 1$. We need an estimate on the number of terms in each of these polynomials.

Both sequences $\tilde{g}^{(j+1)}$ and $\tilde{h}^{(j+1)}$, see (4.1) and (4.2), have components defined in terms of the operator $L(x)$ given by (2.8). In terms of the shifts S^\pm on $\ell^2(\mathbb{Z})$, i.e. $(S^\pm f)_n = f_{n\pm 1}$, one can write

$$(4.21) \quad L(x) = aS^+ + a^-S^- + b.$$

Here a and b are regarded as multiplication operators and a^\pm is the multiplication operator given by $(a^\pm f)_n = a_{n\pm 1}f_n$. An upper bound on the desired number of terms can be obtained by taking both a and b to be constant sequences, i.e., calculating

$$(4.22) \quad \eta^{(j+1)} = \left\langle \delta_n, (S^+ + S^- + I)^{j+1} \delta_n \right\rangle = \sum_{k=0, k \text{ even}}^{j+1} \binom{j+1}{k} \binom{k}{k/2}$$

and

$$(4.23) \quad \xi^{(j+1)} = \left\langle \delta_{n+1}, (S^+ + S^- + I)^{j+1} \delta_n \right\rangle = \sum_{k=0, k \text{ odd}}^{j+1} \binom{j+1}{k} \binom{k}{(k+1)/2}$$

The bounds $\eta^{(j+1)} \leq 2\xi^{(j+1)}$ and $\xi^{(j+1)} \leq 3^j$ readily follow.

Next, using this estimate, we can expand

$$(4.24) \quad a_n(s) \left(\tilde{g}_{n+1}^{(j+1)}(s) - \tilde{g}_n^{(j+1)}(s) \right) = \sum_{k=1}^{2\eta^{(j+1)}} \prod_{\ell=1}^{j+2} d_{n+e_{k,\ell}}$$

and

$$(4.25) \quad \tilde{h}_n^{(j+1)}(s) - \tilde{h}_{n-1}^{(j+1)}(s) = 2 \sum_{k=1}^{2\xi^{(j+1)}} \prod_{\ell=1}^{j+2} d_{n+e'_{k,\ell}}$$

where $d_{n+e} \in \{\pm a_{n+e}(s), \pm b_{n+e}(s)\}$. Here we have inserted a 2 because of the definition of $\tilde{h}^{(j+1)}$, see (4.2). This form enables us to determine the entries of the matrix $D_{n,e}^{(j)}(s)$. In fact, differentiation of the left hand

side of (4.24), with respect to $z \in \{a_m, b_m\}$, yields

$$(4.26) \quad \sum_{k=1}^{2\eta^{(j+1)}} \sum_{o=1}^{j+2} \frac{\partial}{\partial z} d_{n+e_{k,o}} \cdot \prod_{\ell \neq o}^{j+2} d_{n+e_{k,\ell}} = \sum_{|e| \leq \lfloor \frac{j}{2} \rfloor + 1} \frac{\partial}{\partial z} d_{n+e} \sum_{k=1}^{2\eta^{(j+1)}} \sum_{o=1}^{j+2} \delta_e(e_{k,o}) \prod_{\ell \neq o}^{j+2} d_{n+e_{k,\ell}}$$

and an analogous formula holds for the partial derivative of the left hand side of (4.25). These expressions determine the matrix entries of $D_{n,e}^{(j)}(s)$. The bound in (4.15) now follows with

$$(4.27) \quad D_e^{(j)} = \begin{pmatrix} \sum_{k=1}^{2\eta^{(j+1)}} \sum_{o=1}^{j+2} \delta_e(e_{k,o}) & \sum_{k=1}^{2\eta^{(j+1)}} \sum_{o=1}^{j+2} \delta_e(e_{k,o}) \\ 2 \sum_{k=1}^{2\xi^{(j+1)}} \sum_{o=1}^{j+2} \delta_e(e_{k,o}) & 2 \sum_{k=1}^{2\xi^{(j+1)}} \sum_{o=1}^{j+2} \delta_e(e_{k,\xi}) \end{pmatrix}.$$

Summing on e yields,

$$(4.28) \quad \|D^{(j)}\|_\infty = \left\| \sum_{|e| \leq \lfloor \frac{j}{2} \rfloor + 1} D_e^{(j)} \right\|_\infty = 8(j+2)\xi^{(j+1)}$$

where we have used $\eta^{(j+1)} \leq 2\xi^{(j+1)}$. This shows that

$$(4.29) \quad \|D(r)\|_\infty \leq 8 \sum_{j=0}^r |c_{r-j}| \|L(0)\|_2^j (j+2) 3^j,$$

and we are done. \square

We end this section with an analogue of Theorem 2.3 for the hierarchy. Let $\alpha_t^{(r)}$ denote the dynamics corresponding to the Toda hierarchy, i.e., $\alpha_t^{(r)}(A) = A \circ \Phi_t^{(r)}$ for all $A \in \mathcal{A}$.

Theorem 4.4. *Let $r \in \mathbb{N}_0$, $x \in M$, and $\mu > 0$. There exist numbers C_r and v_r for which given any observables $A, B \in \mathcal{A}^{(1)}$, the estimate*

$$(4.30) \quad |\{\alpha_t^{(r)}(A), B\}(x)| \leq C \sum_{n,m \in \mathbb{Z}} \left(\left\| \frac{\partial A}{\partial a_m} \right\|_x + \left\| \frac{\partial A}{\partial b_m} \right\|_x \right) \left(\left\| \frac{\partial B}{\partial a_n} \right\|_x + \left\| \frac{\partial B}{\partial b_n} \right\|_x \right) e^{-\mu(\lceil |n-m|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil - v_r |t|)}$$

holds for all $t \in \mathbb{R}$. Here $4C_r = \|a\|_\infty (1 + e^{\mu/(\lfloor \frac{r}{2} \rfloor + 1)})$ and v_r is as in Theorem 4.2.

Proof. The estimate (4.30) follows as in the proof of Theorem 2.3 using the results of Theorem 4.2 as input. \square

5. RESULTS FOR THE PERTURBED HIERARCHY

The purpose of this section is to demonstrate that the methods from Section 3 also apply to the Toda hierarchy. Many of the proofs follow closely the previous arguments, and so we only sketch the details.

To introduce the relevant class of perturbations, we first recall the Hamiltonian formulation of the hierarchy, see e.g. Section 1.7 of [15]. For each $r \in \mathbb{N}_0$, consider the following formal Hamiltonian

$$(5.1) \quad H_r(x) = \frac{4}{r+2} \sum_{k \in \mathbb{Z}} \sum_{j=0}^r c_{r-j} (\tilde{g}_k^{(j+2)}(x) - \lambda_{r+2}), \quad \lambda_r = \begin{cases} \frac{1}{2^r} \binom{r}{r/2}, & r \text{ even}, \\ 0, & r \text{ odd}, \end{cases}$$

with $\tilde{g}_k^{(j+2)}$ as defined in (4.1).

It can be shown that H_r generates the Toda hierarchy in the sense that for any $A \in \mathcal{A}_0$,

$$(5.2) \quad \frac{d}{dt} \alpha_t^{(r)}(A) = \alpha_t^{(r)}(\{A, H_r\}) = \left\{ \alpha_t^{(r)}(A), H_r \right\}$$

where $\alpha_t^{(r)}$ is the dynamics associated with the Toda hierarchy as introduced before Theorem 4.4. In fact, it is shown in e.g. [15, Thm. 1.71] that

$$(5.3) \quad \frac{\partial H_r}{\partial a_n}(x) = 4g_n(x, r) \quad \text{and} \quad A_n \frac{\partial H_r}{\partial b_n}(x) = 4h_n(x, r)$$

and therefore the evolution equations

$$(5.4) \quad \frac{d}{dt} \alpha_t^{(r)}(A_n) = \alpha_t^{(r)}(\{A_n, H_r\}) = \frac{1}{4} \alpha_t^{(r)} \left(A_n \left(\frac{\partial H_r}{\partial b_{n+1}} - \frac{\partial H_r}{\partial b_n} \right) \right)$$

and similarly,

$$(5.5) \quad \frac{d}{dt} \alpha_t^{(r)}(B_n) = \alpha_t^{(r)}(\{B_n, H_r\}) = \frac{1}{4} \alpha_t^{(r)} \left(A_n \frac{\partial H_r}{\partial a_n} - A_{n-1} \frac{\partial H_r}{\partial a_{n-1}} \right)$$

follow; compare with (4.3) and (4.4).

Note that for $r = 0$

$$(5.6) \quad H_0(x) = \sum_{k \in \mathbb{Z}} (2b_k^2 + 4a_k^2 - 1).$$

is different from (2.5). However, the equations of motion (2.4) are the same as above.

We can now introduce perturbations as in Section 3. Fix $r \in \mathbb{N}_0$ and let $W : \mathbb{R} \rightarrow [0, \infty)$ satisfy $W \in C^2(\mathbb{R})$. Consider the formal Hamiltonian

$$(5.7) \quad H_r^w = H_r + \sum_{n \in \mathbb{Z}} W_n$$

where H_r is as in (5.1) above and W_n is the observable with $W_n(x) = W(\ln(4a_n^2))$.

The equations of motion corresponding to H_r^w are

$$(5.8) \quad \dot{a}_n^w(t, r) = a_n^w(t, r) (g_{n+1}(t, r) - g_n(t, r))$$

and

$$(5.9) \quad \dot{b}_n^w(t, r) = h_n(t, r) - h_{n-1}(t, r) + R_n(t)$$

where

$$(5.10) \quad R_n(t) = \frac{1}{2} [W'(\ln(4a_n^w(t, r)^2)) - W'(\ln(4a_{n-1}^w(t, r)^2))].$$

Again, local existence and uniqueness of solutions of (5.8) and (5.9), corresponding to initial conditions $x \in M_0$, follows from standard results, [1, Thm. 4.1.5]. As before, let us denote by $\Phi_{t,r}^w$ the perturbed flow of the Toda hierarchy, i.e., the function $\Phi_{t,r}^w(x) = \{(a_n^w(t, r), b_n^w(t, r))\}$. Our arguments apply to the set of initial conditions $M_{b,r} = M_{b,r}(W)$ with bounded trajectories, i.e., for which there exists numbers $C_1, C_2 < \infty$ with

$$(5.11) \quad \sup_{t \in \mathbb{R}} \|\Phi_{t,r}^w(x)\|_M \leq C_1 \quad \text{and} \quad \sup_{t \in \mathbb{R}} \sup_{n \in \mathbb{Z}} \frac{1}{|a_n^w(t, r)|} \leq C_2.$$

For initial conditions $x \in M_{b,r}$, there are results similar to the two main estimates from Section 3. We first state an analogue of Theorem 3.1.

Theorem 5.1. *Fix $r \in \mathbb{N}_0$, $W \in C^2(\mathbb{R})$ with $W'' \in L^\infty(\mathbb{R})$, and $x \in M_{b,r}$. For each $\mu > 0$, there exist a number $v_r^w = v_r^w(\mu, x)$ for which given any $n, m \in \mathbb{Z}$, the estimate*

$$(5.12) \quad \max \left[\left| \frac{\partial}{\partial z} a_n^w(t, r) \right|, \left| \frac{\partial}{\partial z} b_n^w(t, r) \right| \right] \leq e^{-\mu(\lceil |n-m|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil - v_r^w |t|)},$$

holds for all $t \in \mathbb{R}$. Here $z \in \{a_m, b_m\}$.

Proof. We will follow closely the proof of Theorem 3.1 using Theorem 4.2 as input. Take $t \geq 0$ and, as before, introduce

$$(5.13) \quad F_n^w(t, r) = \begin{pmatrix} a_n^w(t, r) \\ b_n^w(t, r) \end{pmatrix}.$$

We will suppress the dependence on r . Since W is sufficiently smooth, F_n^w is differentiable with respect to $z \in \{a_m, b_m\}$ and the bound

$$(5.14) \quad \left| \frac{\partial}{\partial z} F_n^w(t) \right| \leq \left| \frac{\partial}{\partial z} F_n^w(0) \right| + \sum_{j=0}^{r+1} \alpha_j \sum_{|e| \leq \beta_j} \int_0^t D_e^{(j)} \left| \frac{\partial}{\partial z} F_{n+e}^w(s) \right| ds.$$

follows as in Theorem 4.2. Here we have denoted by

$$(5.15) \quad \alpha_j = |c_{r-j}| C_1^{j+1} \quad \text{for } 0 \leq j \leq r \quad \text{with } \alpha_{r+1} = C_2 \|W''\|_\infty$$

and

$$(5.16) \quad \beta_j = \lfloor j/2 \rfloor + 1 \quad \text{for } 0 \leq j \leq r \quad \text{with } \beta_{r+1} = 1.$$

Moreover, for $0 \leq j \leq r$, we have taken $D_e^{(j)}$ as in the proof of Theorem 4.2 and set

$$(5.17) \quad D_e^{(r+1)} = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \quad \text{for } e = -1, 0 \quad \text{and} \quad D_1^{(r+1)} = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}.$$

Taking $z = a_m$ and iterating yields

$$(5.18) \quad \left| \frac{\partial}{\partial z} F_n^w(t) \right| \leq \sum_{k=\lceil |n-m|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil}^{\infty} \frac{t^k}{k!} (D_r^w)^k \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

with

$$(5.19) \quad D_r^w = \sum_{j=0}^{r+1} \alpha_j D^{(j)} \quad \text{and} \quad D^{(j)} = \sum_{|e| \leq \beta_j} D_e^{(j)}.$$

A similar estimate holds for the case $z = b_m$.

As in the proof of Theorem 4.2, the bound (5.12) now follows with

$$(5.20) \quad v_r^w = \|D_r^w\|_{\infty} \left(e^{\mu+1} + \frac{1}{\mu} \right).$$

Note also that

$$(5.21) \quad \|D_r^w\|_{\infty} \leq C_1 \sum_{j=0}^r |c_{r-j}| C_1^j \|D^{(j)}\|_{\infty} + 2C_2 \|W''\|_{\infty},$$

and so the estimate from Lemma 4.3 applies here as well. \square

The interpolation argument proven in Section 3 generalizes to the hierarchy as well. To see this, we first prove an analogue of Lemma 3.2.

Lemma 5.2. *Fix $r \in \mathbb{N}_0$, $\mu > 0$, and let $\mathbf{x} \in M_0$. There exists a number $C = C(r, \mu, \mathbf{x}) > 0$ and a function h , depending on r , μ , and \mathbf{x} , for which given any $n, k, \ell \in \mathbb{Z}$, the estimate*

$$(5.22) \quad \max \left[\left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} a_n(t, r) \right|, \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} b_n(t, r) \right| \right] \leq C e^{-\mu(\lceil |n-\ell|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil)} e^{-\mu(\lceil |n-k|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil)} e^{2\mu v_r |t|} h(t)$$

holds for all $t \in \mathbb{R}$. Here $a_n(t, r)$ and $b_n(t, r)$ are the solutions of (2.4) with initial condition $\mathbf{x} \in M$, $\frac{\partial}{\partial \tilde{b}_k} = \frac{\partial}{\partial b_{k+1}} - \frac{\partial}{\partial b_k}$, $z \in \{a_{\ell}, b_{\ell}\}$, and the number v_r is as in Theorem 2.1. The function h grows at most exponentially.

Proof. Again, since this proof follows closely the arguments of Lemma 3.2, we will only sketch the details.

Keeping with the previous notation, it is clear that

$$(5.23) \quad \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) = \sum_{j=0}^r c_{r-j} \sum_{|e| \leq \lfloor j/2 \rfloor + 1} \int_0^t \frac{\partial}{\partial z} \left(D_{n,e}^{(j)}(s) \frac{\partial}{\partial \tilde{b}_k} F_{n+e}(s) \right) ds.$$

Differentiation with respect to z produces two terms. We begin by estimating the first term, i.e., the one that contains $\frac{\partial}{\partial z} D_{n,e}^{(j)}(s)$. Using Theorem 4.2, it is clear that if $|e| \leq \lfloor j/2 \rfloor + 1$, then

$$(5.24) \quad \left| \frac{\partial}{\partial \tilde{b}_k} F_{n+e}(s) \right| \leq e^{\mu} e^{-\mu(\lceil |n-k|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil - v_r |s|)} \begin{pmatrix} 1 \\ 1 \end{pmatrix}.$$

For convenience, let us denote by

$$(5.25) \quad E(x) = e^{-\mu(\lceil |x|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil)}.$$

Using (5.24), one can show that

$$(5.26) \quad \sum_{|e| \leq \lfloor j/2 \rfloor + 1} \int_0^{|t|} \left| \frac{\partial}{\partial z} D_{n,e}^{(j)}(s) \frac{\partial}{\partial \tilde{b}_k} F_{n+e}(s) \right| ds \leq \gamma_j(\mu) E(n-k) E(n-\ell) \int_0^{|t|} e^{2\mu v_r s} ds \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

where

$$(5.27) \quad \gamma_j(\mu) = 8e^{2\mu} \|L(0)\|_2^j (j+1)(j+2).$$

This proves that

$$(5.28) \quad \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) \right| \leq C_\mu(r) E(n-k) E(n-\ell) \int_0^{|t|} e^{2\mu v_r s} ds \begin{pmatrix} 1 \\ 1 \end{pmatrix} \\ + \|L(0)\|_2 \sum_{j=0}^r |c_{r-j}| \|L(0)\|_2^j \sum_{|e| \leq \lfloor j/2 \rfloor + 1} \int_0^{|t|} D_e^{(j)} \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_{n+e}(s) \right| ds$$

where

$$(5.29) \quad C_\mu(r) = \sum_{j=0}^r |c_{r-j}| \gamma_j(\mu),$$

and $D_e^{(j)}$ is as in (4.15). Iteration, as in the proof of Lemma 3.2, yields

$$(5.30) \quad \left| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) \right| \leq \frac{C_\mu(r)}{2\mu v_r} E(n-k) E(n-\ell) \sum_{m=0}^{\infty} \left(\frac{\|L(0)\|_2}{2\mu v_r} e^{2\mu} \right)^m \sum_{p=m+1}^{\infty} \frac{(2\mu v_r |t|)^p}{p!} \tilde{D}(r)^m \begin{pmatrix} 1 \\ 1 \end{pmatrix},$$

where

$$(5.31) \quad \tilde{D}(r) = \sum_{j=0}^r |c_{r-j}| \|L(0)\|_2^j \sum_{|e| \leq \lfloor j/2 \rfloor + 1} D_e^{(j)}.$$

Taking ∞ -norms, and using (3.34), we see that

$$(5.32) \quad \left\| \frac{\partial^2}{\partial z \partial \tilde{b}_k} F_n(t) \right\|_{\infty} \leq C E(n-k) E(n-\ell) e^{2\mu v_r |t|} h(t)$$

where

$$(5.33) \quad C = \frac{C_\mu(r)}{2\mu v_r} \quad \text{and with } \beta = \frac{e^{2\mu} \|L(0)\|_2 \tilde{D}(r)}{2\mu v_r} - 1$$

h is as in (3.36). This proves (5.22) and we are done. \square

Given Lemma 5.2 above, the analogue of Theorem 3.3 follows with only minor modifications. We state it below.

Theorem 5.3. Fix $r \in \mathbb{N}_0$, $W \in C^2(\mathbb{R})$ with $W', W'' \in L^\infty(\mathbb{R})$, and let $x \in M_{b,r}$. For each $\mu > 0$ and any $\varepsilon > 0$, there are positive numbers $C = C(r, \varepsilon)$, $D = D(r, \varepsilon, \mu, x, W)$, and $\delta = \delta(r, \varepsilon, \mu, x, W)$ such that

$$(5.34) \quad \max \left[\left\| \frac{\partial}{\partial z} a_n^w(t, r) \right\|, \left\| \frac{\partial}{\partial z} b_n^w(t, r) \right\| \right] \leq C G_{\mu_r}(|n-m|) e^{(\mu+\varepsilon)v_r|t|} \left[1 + D \left(e^{\delta|t|} - 1 \right) \right]$$

holds for all $t \in \mathbb{R}$. Here $v_r = v_r(\mu + \varepsilon, x)$ is as in Theorem 4.2, $\mu_r = \mu / (\lfloor \frac{r}{2} \rfloor + 1)$, and $z \in \{a_m, b_m\}$.

Proof. Interpolating as before, it is clear that for any $x \in M$, the bound

$$(5.35) \quad \left\| \left[\frac{\partial}{\partial z} \tilde{F}_n^w(t) \right] (x) \right\|_{\infty} \leq \left\| \left[\frac{\partial}{\partial z} \tilde{F}_n(t) \right] (x) \right\|_{\infty} + \sum_{k \in \mathbb{Z}} \int_0^{|t|} \left\| \left[D_k(s; n) \frac{\partial}{\partial z} \tilde{F}_k^w(s) \right] (x) \right\|_{\infty} ds \\ + \frac{\|W'\|_{\infty}}{2} \sum_{k, \ell \in \mathbb{Z}} \int_0^{|t|} \left\| \left[D_{k, \ell}(s; n) \frac{\partial}{\partial z} \tilde{F}_\ell^w(s) \right] (x) \right\|_{\infty} ds$$

follows as in the proof of (3.55); here all quantities depend now also on r , but we have suppressed this in our notation.

With Theorem 4.2, it is clear that

$$(5.36) \quad \left\| \left[\frac{\partial}{\partial z} \tilde{F}_n(t) \right] (x) \right\|_{\infty} = \left\| \frac{\partial}{\partial z} F_n(t) \right\|_{\infty} \\ \leq e^{-(\mu+\varepsilon)(\lceil |n-m| / (\lfloor \frac{r}{2} \rfloor + 1) \rceil - v_r|t|)} \\ \leq C_{\varepsilon, r} G_{\mu_r}(|n-m|) e^{(\mu+\varepsilon)v_r|t|},$$

where we have set

$$(5.37) \quad C_{\varepsilon,r} = \sup_{x \geq 0} (1+x)^2 e^{-\varepsilon(\lceil x/(\lfloor \frac{r}{2} \rfloor + 1) \rceil)}, \quad \mu_r = \frac{\mu}{\lfloor \frac{r}{2} \rfloor + 1}, \quad \text{and} \quad v_r = v_r(x, \mu + \varepsilon).$$

Similarly, for each $x \in M_{b,r}$, the matrix appearing in the second term satisfies

$$(5.38) \quad \begin{aligned} \|[D_k(s; n)](x)\|_\infty &\leq C_2 \|W''\|_\infty e^{\mu+\varepsilon} e^{-(\mu+\varepsilon)(\lceil |n-k|/(\lfloor \frac{r}{2} \rfloor + 1) \rceil - v_r(s)(|t|-s))} \\ &\leq C_2 \|W''\|_\infty e^{\mu+\varepsilon} C_{\varepsilon,r} G_{\mu_r}(|n-k|) e^{(\mu+\varepsilon)v_r^*(|t|-s)} \end{aligned}$$

with

$$(5.39) \quad v_r(s) = v_r(\Phi_{s,r}^w(x), \mu + \varepsilon) \quad \text{and} \quad v_r^* = \sup_{s \in \mathbb{R}} v_r(s).$$

Lastly, the bound

$$(5.40) \quad \|[D_{k,\ell}(s; n)](x)\|_\infty \leq \tilde{C} G_{\mu_r}(|n-\ell|) G_{\mu_r}(|n-k|) e^{2(\gamma+1)(\mu+\varepsilon)v_r^*(|t|-s)}$$

follows as in the proof of (3.63); with (possibly) different values of \tilde{C} and γ .

Iteration yields (5.34) as in the proof of Theorem 3.3. \square

6. LOCALITY BOUNDS FOR MORE GENERAL INITIAL CONDITIONS

In this section, we return to the class of perturbations considered in Section 3, see also Section 5. The goal here is to prove a locality estimate for more general initial conditions. Let us recall the basic set-up.

Fix $W : \mathbb{R} \rightarrow [0, \infty)$ satisfying $W \in C^2(\mathbb{R})$ with $W', W'' \in L^\infty(\mathbb{R})$. The formal Hamiltonian is given by

$$(6.1) \quad H^w = H + \sum_{n \in \mathbb{Z}} W_n,$$

where H is the Toda Hamiltonian as in (2.5), and for each n , the perturbation W_n is taken to be the observable $W_n(x) = W(\ln(4a_n^2))$. The corresponding equations of motion are

$$(6.2) \quad \begin{aligned} \dot{a}_n^w(t) &= a_n^w(t) (b_{n+1}^w(t) - b_n^w(t)) \\ \dot{b}_n^w(t) &= 2(a_n^w(t)^2 - a_{n-1}^w(t)^2) + R_n(t) \end{aligned}$$

where

$$(6.3) \quad R_n(t) = \frac{1}{2} [W'(\ln(4a_n^w(t)^2)) - W'(\ln(4a_{n-1}^w(t)^2))].$$

As we discussed before, local existence and uniqueness of solutions of (6.2) corresponding to initial conditions $x \in M_0$ follows again from Theorem 4.1.5 in [1]. For $x \in M_0$, let us again denote by $\Phi_t^w(x) = \{(a_n^w(t), b_n^w(t))\}$ the perturbed Toda flow. If we set, $L(t) = L(\Phi_t^w(x))$ and $P(t) = P(\Phi_t^w(x))$ with L and P as in (2.8) and (2.9), it is easy to check that

$$(6.4) \quad \frac{d}{dt} L(t) = [P(t), L(t)] + R(t),$$

where the multiplication operator $R(t) : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$ is defined by

$$(6.5) \quad [R(t)f]_n = R_n(t)f_n.$$

From this definition, it is clear that

$$(6.6) \quad \|R(t)\|_2 \leq \|W'\|_\infty,$$

and hence the operator norm of $R(t)$ is bounded uniformly in both t and the initial condition $x \in M_0$.

One can check that $P(t)$ still corresponds to a family of unitary propagators, which we denote by $U(t, s)$. A short calculation shows then that $\tilde{L}(t) = U(t, s)^* L(t) U(t, s)$ satisfies

$$(6.7) \quad \frac{d}{dt} \tilde{L}(t) = U(t, s)^* R(t) U(t, s).$$

In this case,

$$(6.8) \quad \|L(t)\|_2 = \|\tilde{L}(t)\|_2 \leq \|\tilde{L}(0)\|_2 + \int_0^{|t|} \|R(s)\|_2 ds \leq \|L(0)\|_2 + \|W'\|_\infty |t|.$$

Since the bound

$$(6.9) \quad \max(\sup_n |a_n^w(t)|, \sup_n |b_n^w(t)|) \leq \|L(t)\|_2$$

holds, (6.8) produces a linear bound on the growth of solutions. From this, the existence of global solutions follows from Proposition 4.1.22 in [1].

A result analogous to Theorem 2.1 follows.

Theorem 6.1. *Let $x = \{(a_n, b_n)\}_{n \in \mathbb{Z}} \in M_0$ with $a_* = \inf_n |a_n| > 0$. Then for any $\mu > 0$ and $n, m \in \mathbb{Z}$ the estimate*

$$(6.10) \quad \max \left[\left| \frac{\partial}{\partial z} \ln(a_n^w(t)^2) \right|, \left| \frac{\partial}{\partial z} b_n^w(t) \right| \right] \leq \max \left(1, \frac{2}{a_*} \right) e^{-\mu(|n-m|-v^w(t))}$$

holds for all $t \in \mathbb{R}$. Here $z \in \{a_m, b_m\}$ and $v^w(t)$ is an explicit cubic polynomial with $v^w(0) = 0$.

Proof. Fix $x \in M_0$ with $a_* > 0$. Global existence on M guarantees that for each $n \in \mathbb{Z}$, the function $F_n^w : \mathbb{R} \rightarrow \mathbb{R}^2$ given by

$$(6.11) \quad F_n^w(t) = \begin{pmatrix} \ln(a_n^w(t)^2) \\ b_n^w(t) \end{pmatrix}$$

is well-defined for all $t \in \mathbb{R}$. In fact, since $a_* > 0$, the equations of motion ensure that

$$(6.12) \quad a_n^w(t) = a_n^w(0) \exp \left(\int_0^t (b_{n+1}^w(s) - b_n^w(s)) ds \right)$$

preventing a singularity in the logarithm. It is then clear that (6.2) implies

$$(6.13) \quad F_n^w(t) = F_n^w(0) + 2 \int_0^t \left(a_n^w(s)^2 - a_{n-1}^w(s)^2 + \frac{1}{2} R_n(s) \right) ds.$$

For any $z \in \{a_m, b_m\}$, a relation similar to (2.18) holds, i.e.

$$(6.14) \quad \frac{\partial}{\partial z} F_n^w(t) = \frac{\partial}{\partial z} F_n^w(0) + 2 \sum_{|e| \leq 1} \int_0^t D_{n,e}^w(s) \frac{\partial}{\partial z} F_{n+e}^w(s) ds,$$

where

$$(6.15) \quad D_{n,e}^w(s) = \begin{pmatrix} 0 & \delta_1(e) - \delta_0(e) \\ \left(a_n^w(s)^2 + \frac{\|W''\|_\infty}{4} \right) \delta_0(e) - \left(a_{n-1}^w(s)^2 + \frac{\|W''\|_\infty}{4} \right) \delta_{-1}(e) & 0 \end{pmatrix}$$

The bounds in (6.8) and (6.9) show that

$$(6.16) \quad a_n^w(s)^2 \leq (\|L(0)\|_2 + \|W'\|_\infty |s|)^2,$$

and so we find that

$$(6.17) \quad \left| \frac{\partial}{\partial z} F_n^w(t) \right| \leq \left| \frac{\partial}{\partial z} F_n^w(0) \right| + \sum_{|e| \leq 1} \int_0^{|t|} g(s) D_e^w \left| \frac{\partial}{\partial z} F_{n+e}^w(s) \right| ds$$

holds for any $t \in \mathbb{R}$. Here we have denoted by

$$(6.18) \quad g(s) = 1 + (\|L(0)\|_2 + \|W'\|_\infty |s|)^2 + \frac{\|W''\|_\infty}{4} \quad \text{and} \quad D_e^w = \begin{pmatrix} 0 & \delta_1(e) + \delta_0(e) \\ \delta_0(e) + \delta_{-1}(e) & 0 \end{pmatrix}$$

Iteration (with $z = a_m$) yields

$$(6.19) \quad \left| \frac{\partial}{\partial z} F_n^w(t) \right| \leq \sum_{k=|n-m|}^{\infty} \sum_{|e_1| \leq 1} \cdots \sum_{|e_k| \leq 1} \int_0^{|t|} \int_0^{t_1} \cdots \int_0^{t_{k-1}} g(t_1) \cdots g(t_k)$$

$$(6.20) \quad \times D_{e_1}^w \cdots D_{e_k}^w \left| \frac{\partial}{\partial z} F_{n+e_1+\cdots+e_k}^w(0) \right| dt_k \cdots dt_1$$

$$(6.21) \quad \leq \frac{2}{|a_m|} \sum_{k=|n-m|}^{\infty} \frac{h(t)^k}{k!} (D^w)^k \begin{pmatrix} 1 \\ 0 \end{pmatrix},$$

where $h(t) = 2 \int_0^{|t|} g(s) ds$ and

$$(6.22) \quad D^w = \frac{1}{2} \sum_{|e| \leq 1} D_e^w = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$

In the case that $z = b_m$, it is clear that

$$(6.23) \quad \left| \frac{\partial}{\partial z} F_n^w(t) \right| \leq \sum_{k=|n-m|}^{\infty} \frac{h(t)^k}{k!} (D^w)^k \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Taking the infinity norm, we get for any $\mu > 0$

$$(6.24) \quad \left\| \frac{\partial}{\partial z} F^w(t) \right\|_{\infty} \leq \max \left(1, \frac{2}{|a_m|} \right) \sum_{k=|n-m|}^{\infty} \frac{h(t)^k}{k!} \\ \leq \max \left(1, \frac{2}{a_*} \right) e^{-\mu(|n-m|-v^w(t))},$$

where $v^w(t) = h(t) (e^{\mu+1} + \mu^{-1})$. □

Again, everything extends to the hierarchy. Since the statements are clear, we do not rewrite them for the sake of brevity.

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APPENDIX A. EXISTENCE OF BOUNDED SOLUTIONS FOR HAMILTONIAN SYSTEMS

In this appendix we want to look at a general Hamiltonian system with nearest neighbor interaction:

$$(A.1) \quad H(x) = \sum_{n \in \mathbb{Z}} \left(\frac{p_n^2}{2} + V(q_{n+1} - q_n) \right),$$

where $x = \{(p_n, q_n)\}_{n \in \mathbb{Z}}$. Let $V \in C^2(\mathbb{R})$ with $V(x) \geq 0$ and $V(0) = V'(0) = 0$ such that 0 is a fixed point of the system. Since we want to obtain bounded solutions we will assume that our interaction potential is confining in the sense that $V(x) \rightarrow +\infty$ as $x \rightarrow \pm\infty$. Since the uniform motion $q_n(t) = q^0 + p^0 t$ (with q^0, p^0 some real constants) of the system is unbounded we switch to relative coordinates $r_n = q_{n+1} - q_n$ in which the equation of motions read

$$(A.2) \quad \dot{r}_n = p_{n+1} - p_n, \quad \dot{p}_n = V'(r_n) - V'(r_{n-1}).$$

We will consider this system in the Hilbert space $\mathcal{X} = \ell^2(\mathbb{Z}) \times \ell^2(\mathbb{Z})$.

Theorem A.1. *Suppose $V \in C^2(\mathbb{R})$ such that 0 is a unique global minimum with $V(0) = V'(0) = 0$, $V''(0) > 0$ and $V(x) \rightarrow +\infty$ as $|x| \rightarrow \infty$.*

Then the system (A.2) has a unique global solution in \mathcal{X} for which the energy

$$(A.3) \quad H(p, r) = \sum_{n \in \mathbb{Z}} \left(\frac{p_n^2}{2} + V(r_n) \right)$$

is finite and conserved. This solution is C^1 with respect to the initial condition. Moreover, 0 is a stable fixed point and all solutions satisfy $\|(p(t), r(t))\|_2 \leq C$ as well as $\|(p(t), r(t))\|_{\infty} \leq C$, where the constant C depends only on the initial condition.

Proof. First of all note that by our assumption on V we can find constants c_R and C_R for every $R > 0$ such that $|V(x)| \leq C_R x^2$, $|V'(x)| \leq C_R |x|$ and $|V(x)| \geq c_R x^2$ for $|x| \leq R$.

In particular, for $r \in \ell^2(\mathbb{Z})$ with $\|r\|_2 \leq R$ we have $\|V'(r)\|_2 \leq C_R \|r\|_2$ and it follows that the map $r_n \mapsto V'(r_n)$ is C^1 on $\ell^2(\mathbb{Z})$. Since the same is true for the shift operator $x_n \mapsto x_{n-1}$, our vector field is C^1 and local existence and uniqueness follow from standard results [1, Thm. 4.1.5]. This also implies that the flow is C^1 with respect to the initial condition [1, Lem. 4.1.9]. Moreover, $|H(p, r)| \leq C_R \|(p, r)\|_2$ implies that H is finite on \mathcal{X} and a straightforward calculation shows that it is conserved by the flow.

Moreover, $H(p, r) \geq c_1 \varepsilon$ for $\|(p, r)\|_2 = \varepsilon$ and $\varepsilon \leq 1$. Hence [1, Thm. 4.3.11] shows that 0 is a stable fixed point.

Finally, if $V(x) \rightarrow +\infty$ there is a constant M_E such that $|x| \leq M_E$ whenever $|V(x)| \leq E$. Hence setting $E = H(p(0), r(0))$ we have $H(p(t), r(t)) = E$ implying $\|p(t)\|_2 \leq \sqrt{2E}$ and $\|r(t)\|_\infty \leq M_E$. But this implies $\|r(t)\|_2 \leq c_{M_E}^{-1} \|V(r)\|_2 \leq c_{M_E}^{-1} \sqrt{E}$. Hence our vector field remains bounded along integral curves on finite t intervals and hence all solutions are global in time by [1, Prop. 4.1.22]. \square

Note that, using $q_n(t) = q_n(0) + \int_0^t p_n(s) ds$, for our original variables we get

$$(A.4) \quad \|q(t)\|_2 \leq \|q(0)\|_2 + C|t|, \quad \|q(t)\|_\infty \leq \|q(0)\|_\infty + C|t|.$$

Moreover, clearly the Toda potential $V(r) = e^{-r} + r - 1$ satisfies the above assumptions.

Theorem A.2. *Let $x = (p, r) \in \mathcal{X}$ and $\mu > 0$. There exists a number $v = v(\mu, x)$ for which given any $n, m \in \mathbb{Z}$, the bound*

$$(A.5) \quad \max \left[\left| \frac{\partial}{\partial z} p_n(t) \right|, \left| \frac{\partial}{\partial z} r_n(t) \right| \right] \leq C e^{-\mu(|n-m|-v|t|)},$$

holds for all $t \in \mathbb{R}$ and each $z \in \{p_m, r_m\}$, where

$$(A.6) \quad C = C(x) = \max \left(\sup_{(t,n) \in \mathbb{R} \times \mathbb{Z}} |V'(r_n(t))|^{1/2}, 1 \right).$$

In fact, one may take

$$(A.7) \quad v = 2C \left(e^{\mu+1} + \frac{1}{\mu} \right).$$

Proof. Fix $x \in \mathcal{X}$. Our previous theorem guarantees that for each $x \in M$ and $n \in \mathbb{Z}$, the function $F_n : \mathbb{R} \rightarrow \mathbb{R}^2$ given by

$$(A.8) \quad F_n(t; x) = \begin{pmatrix} r_n(t) \\ p_n(t) \end{pmatrix}$$

is well-defined and differentiable with respect to each $z \in \{p_m, r_m\}$. When convenient, we will suppress the dependence of F_n on x . Using the equations of motion, i.e. (A.2), it is clear that

$$(A.9) \quad F_n(t) = F_n(0) + \int_0^t \begin{pmatrix} p_{n+1}(s) - p_n(s) \\ V'(r_n(s)) - V'(r_{n-1}(s)) \end{pmatrix} ds.$$

Differentiating with respect to z we obtain

$$(A.10) \quad \frac{\partial}{\partial z} F_n(t) = \frac{\partial}{\partial z} F_n(0) + \sum_{|e| \leq 1} \int_0^t D_{n,e}(s) \frac{\partial}{\partial z} F_{n+e}(s) ds,$$

with

$$(A.11) \quad D_{n,e}(s) = \begin{pmatrix} 0 & \delta_1(e) - \delta_0(e) \\ V'(r_n)\delta_0(e) - V'(r_{n-1})\delta_{-1}(e) & 0 \end{pmatrix}.$$

Hence

$$(A.12) \quad \left| \frac{\partial}{\partial z} F_n(t) \right| \leq \left| \frac{\partial}{\partial z} F_n(0) \right| + \sum_{|e| \leq 1} \int_0^{|t|} D_e \left| \frac{\partial}{\partial z} F_{n+e}(s) \right| ds,$$

where

$$(A.13) \quad D_e = \begin{pmatrix} 0 & \delta_1(e) + \delta_0(e) \\ C^2(\delta_0(e) + \delta_{-1}(e)) & 0 \end{pmatrix}.$$

Let us now consider the case that $z = r_m$, i.e.,

$$(A.14) \quad \frac{\partial}{\partial r_m} F_n(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \delta_m(n).$$

In this case, iteration yields

$$(A.15) \quad \left| \frac{\partial}{\partial r_m} F_n(t) \right| \leq \sum_{k=0}^{\infty} \frac{|t|^k}{k!} \sum_{|e_1| \leq 1} \cdots \sum_{|e_k| \leq 1} \delta_{m+e_1+\dots+e_k}(n) D_{e_1} \cdots D_{e_k} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\ \leq \sum_{k=|n-m|}^{\infty} \frac{|2Ct|^k}{k!} D^k \begin{pmatrix} 1 \\ 0 \end{pmatrix},$$

where we have set

$$(A.16) \quad D = \frac{1}{2C} \sum_{|e| \leq 1} D_e = \begin{pmatrix} 0 & C^{-1} \\ C & 0 \end{pmatrix}.$$

Again, convergence is guaranteed since $\frac{\partial}{\partial z} F_n(t)$ is continuous and thus bounded on compact time intervals. Using $D^2 = \mathbb{1}$, one obtains

$$(A.17) \quad \left\| \frac{\partial}{\partial r_m} F_n(t) \right\|_{\infty} \leq C \sum_{k=|n-m|}^{\infty} \frac{|2Ct|^k}{k!}$$

The rest follows as in Theorem 2.1. □

Note that in Flaschka variables (2.3) the equations of motion read

$$(A.18) \quad \dot{a}_n(t) = a_n(t) (b_{n+1}(t) - b_n(t)) \quad \text{and} \quad \dot{b}_n(t) = -\frac{1}{2} (V'(-\ln(4a_n^2)) - V'(-\ln(4a_{n-1}^2))) ,$$

and (p, r) will be bounded if and only if (a, a^{-1}, b) are bounded. The fixed point in these new coordinates is $(a_0, b_0) = (\frac{1}{2}, 0)$ and $(p, r) \in \mathcal{X}$ if and only if $(a, b) - (a_0, b_0) = (a - \frac{1}{2}, b) \in \mathcal{X}$.

Finally, by

$$(A.19) \quad a_n(t) = a_n(0) \exp \left(\int_0^t (b_{n+1}(s) - b_n(s)) ds \right)$$

the sign of a_n is preserved under the flow.

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